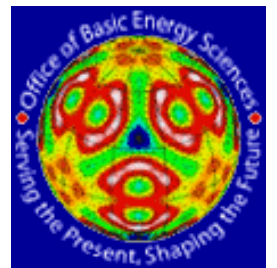


# Studying competition, disorder and nanoscale fluctuations in strongly correlated systems and the necessity of combining methods such as muons, x-rays and neutrons

S.J.L. Billinge

*Department of Applied Physics and Applied Mathematics  
Columbia University,  
CMPMS, Brookhaven National Laboratory*

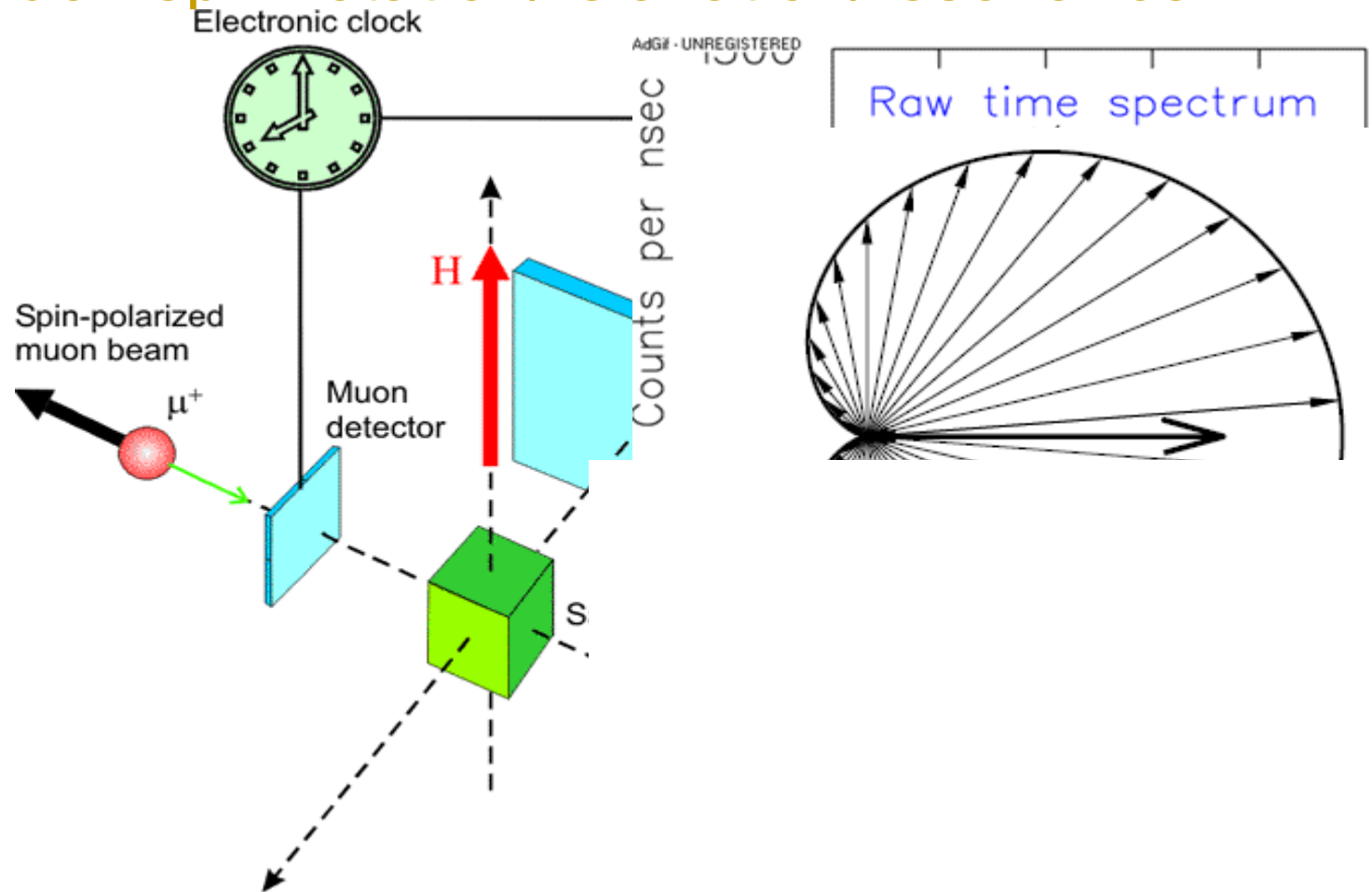


## Why do we love Muons?

- They are a **LOCAL** probe of the magnetic field in a material:
  - Sensitive to all kinds of spin order (or lack of it): antiferromagnets, spin-gap systems, spin glasses, ferromagnets...
  - follow an order parameter as a function of temperature
  - provides information about internal magnetic field distributions (inhomogeneous fields of phase separation)
  - Provides information about spin dynamics and relaxations
  - it works very well at milli-Kelvin temperatures (the incident muons easily pass through the dilution refrigerator windows)

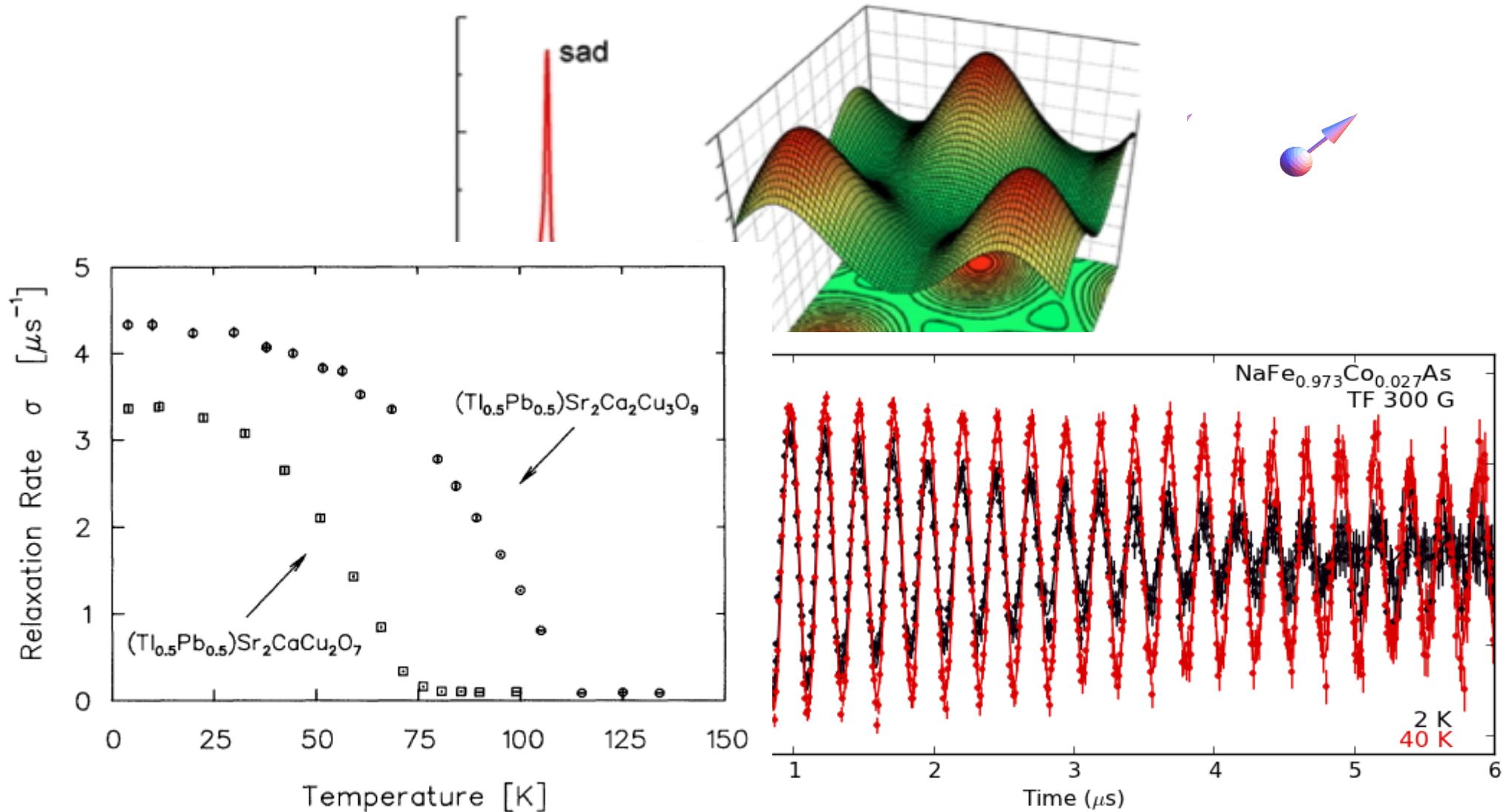
A very versatile probe of magnetism in complex materials

# MuSR: Muon spin rotation/relaxation/resonance

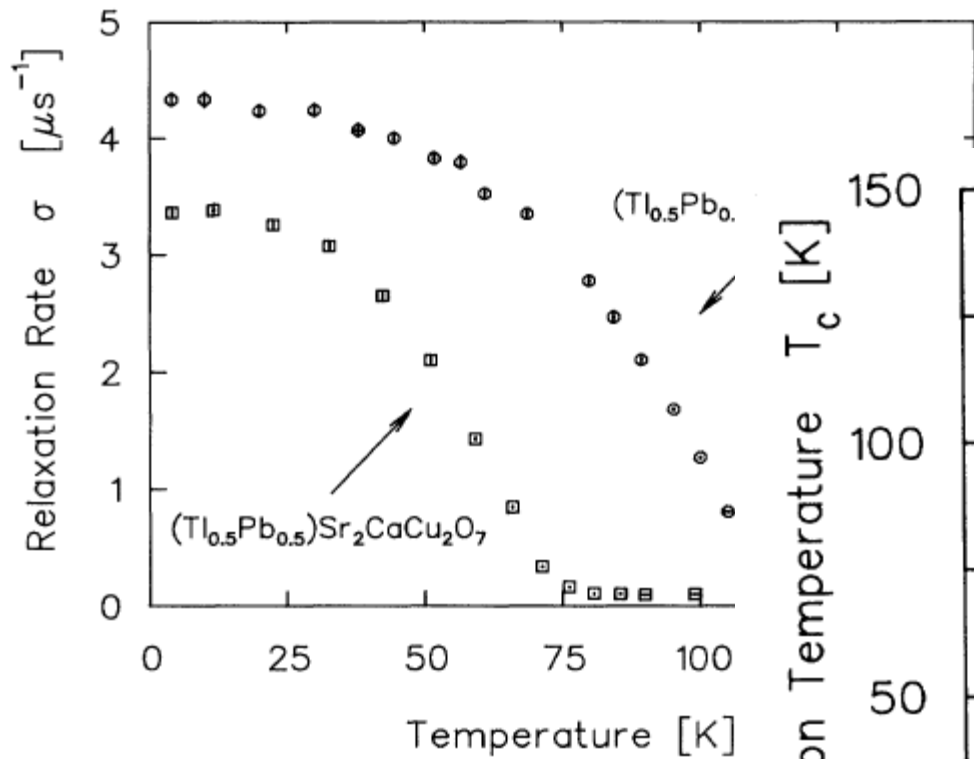


A. Savici, [http://neutron.magnet.fsu.edu/muon\\_relax.html](http://neutron.magnet.fsu.edu/muon_relax.html)

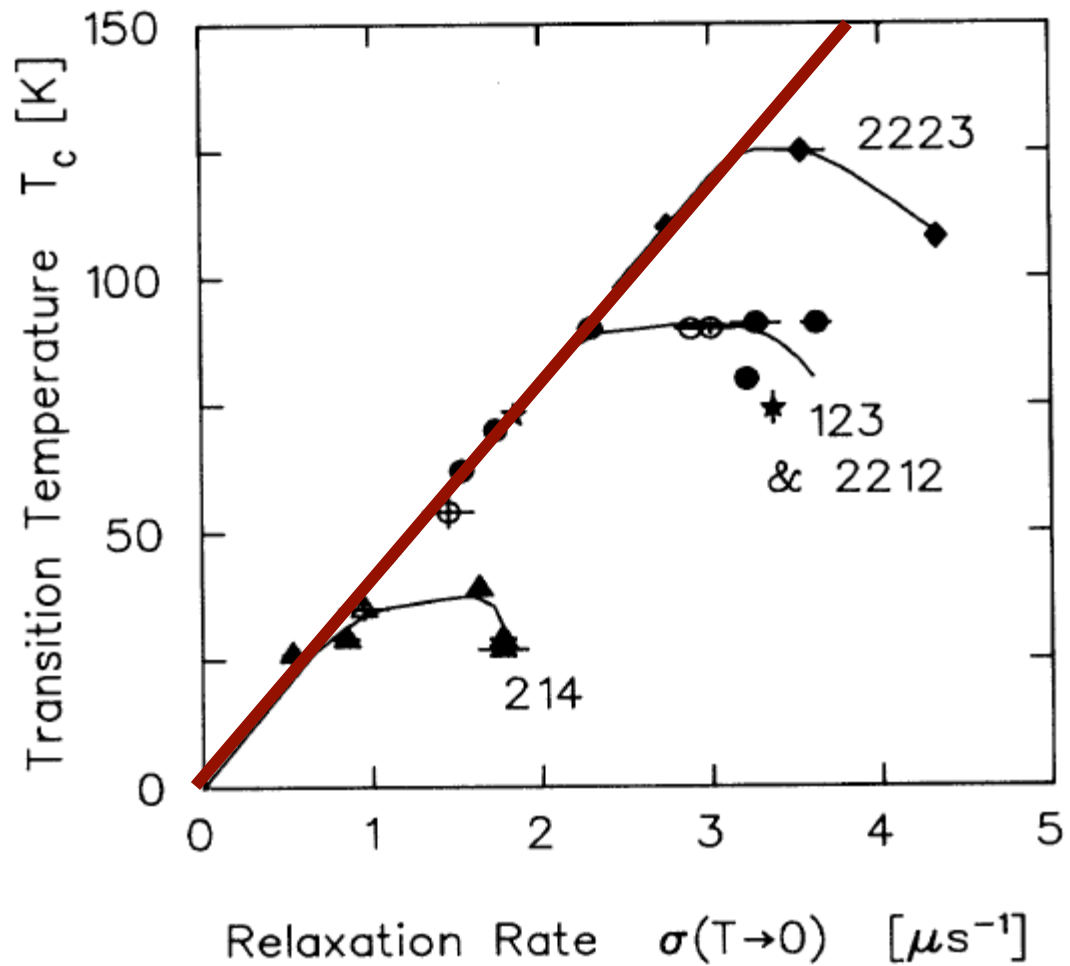
# Inhomogeneous magnetism probed by $\mu$ SR: e.g., superconductivity







- Uemura et al., PRL 1988 => "Uemura Plot"



# Why do we love Muons

- They are a **LOCAL** probe of the magnetic field in a material:
  - Highly complementary to other probes of local atomic and electronic structure from scattering and STM

A very versatile probe of magnetism in complex materials

# Why are muon experiments not more ubiquitous

- ~~• Very difficult to do the experiments?~~
- ~~• Very difficult to analyze the data?~~
- ~~• Many different and easier ways of getting the same information?~~
- Scarcity of muons!
  - CMMS continuous source at TRIUMF in Vancouver, Canada
  - ISIS Pulsed Muon Facility, UK
  - Paul Scherrer Institute Muon Facility, Switzerland
  - RIKEN-RAL Muon Facility, Japan

# Complex materials

- Photovoltaics with improved efficiency
  - Nanoparticles in the light collecting layer
- High energy density batteries and fuel cells
  - Electrodes
  - Electrolytes
  - Catalysts
- Electrical transport
  - High temperature superconductors
- Next generation electronics
  - Spintronics
  - Multiferroics

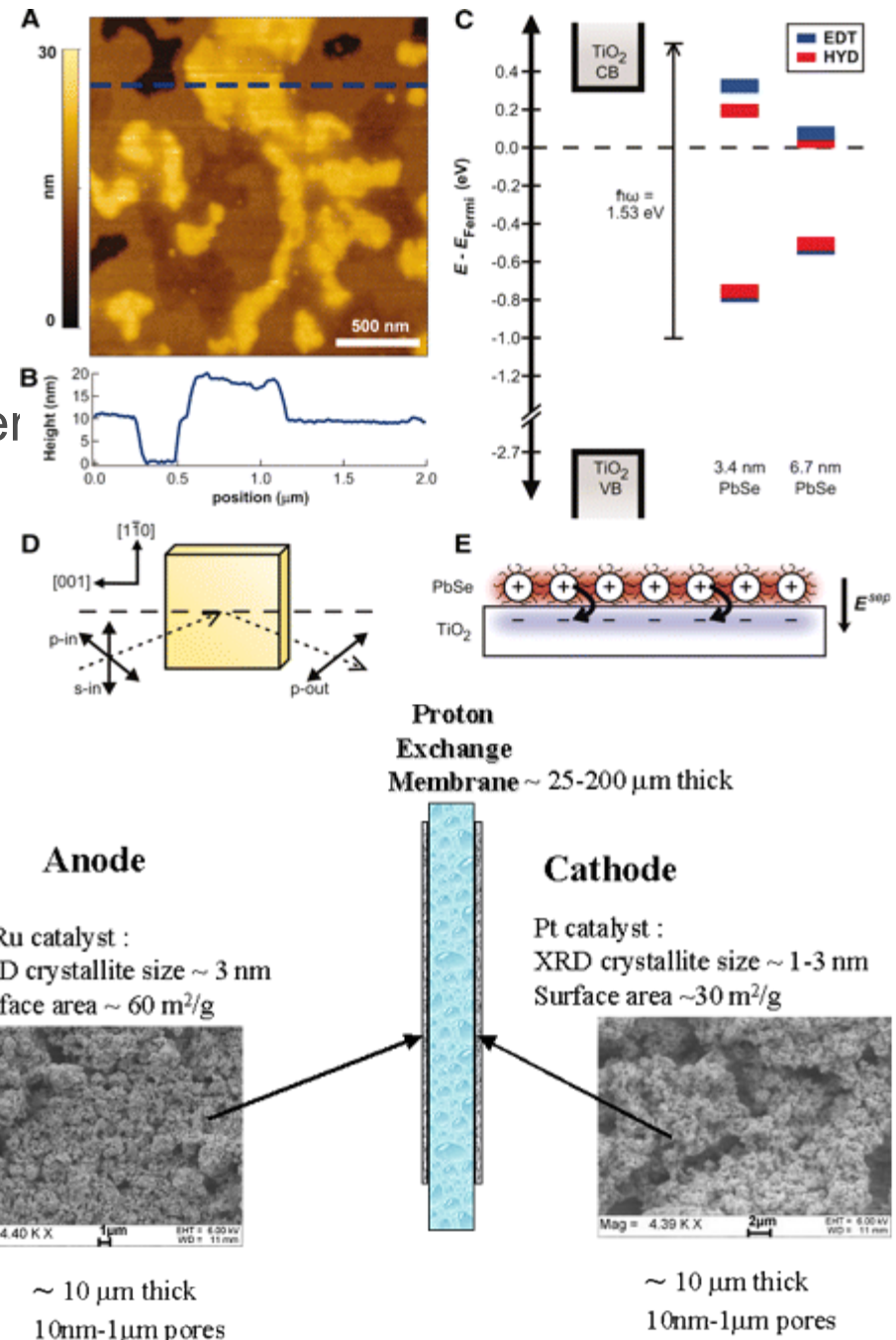
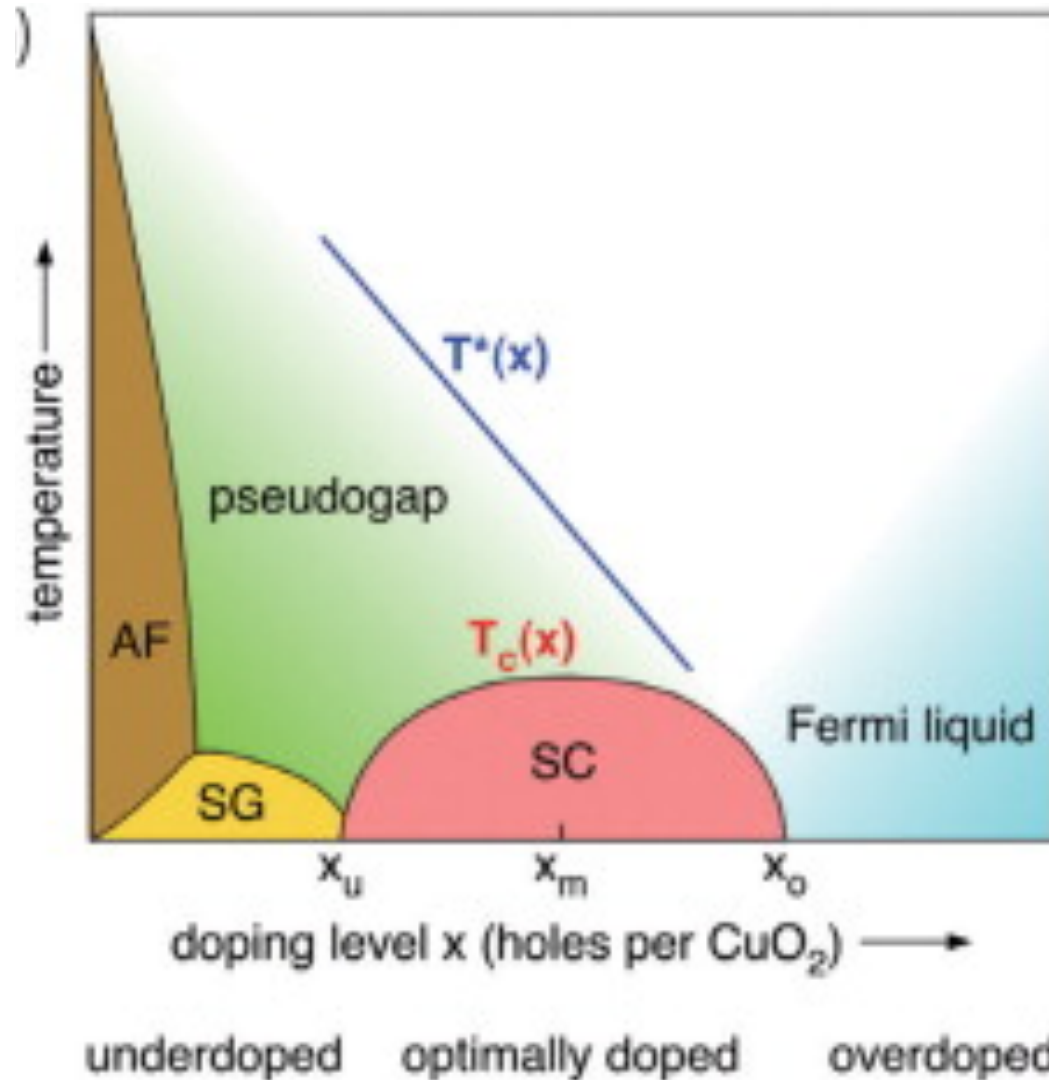


Image credits: 10.1126/science.1185509  
 U. Uppsala

# Materials with colossal emergent properties close to a metal-insulator transition

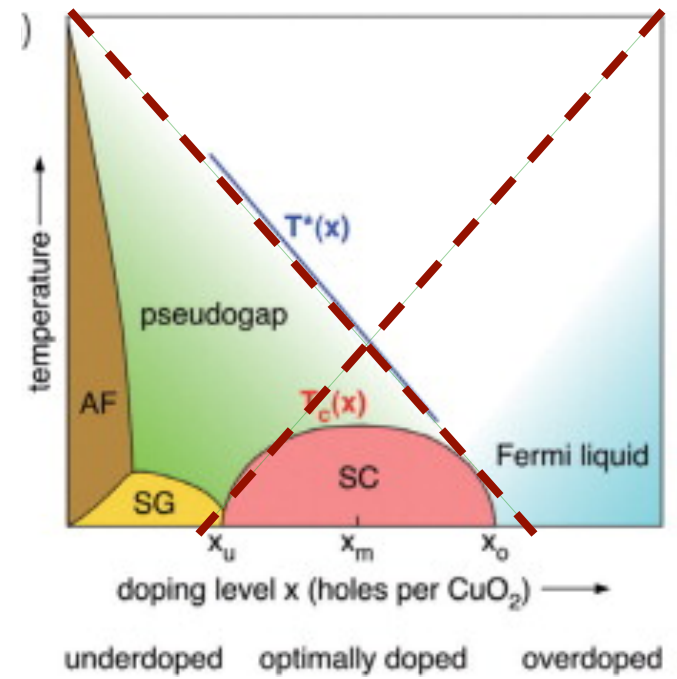
# Materials with colossal emergent properties close to a metal-insulator transition



Generic high temperature superconducting cuprate phase diagram  
Annette Bussman-Holder and K. A. Muller 2008 review

# Cuprate Superconductivity

- Basic picture of doped, half-filled Mott insulator
- Antiferromagnetism and spin fluctuations
- Strange metal
- Spin gap  $\rightarrow$  pseudogap
- Is it an X or a V?

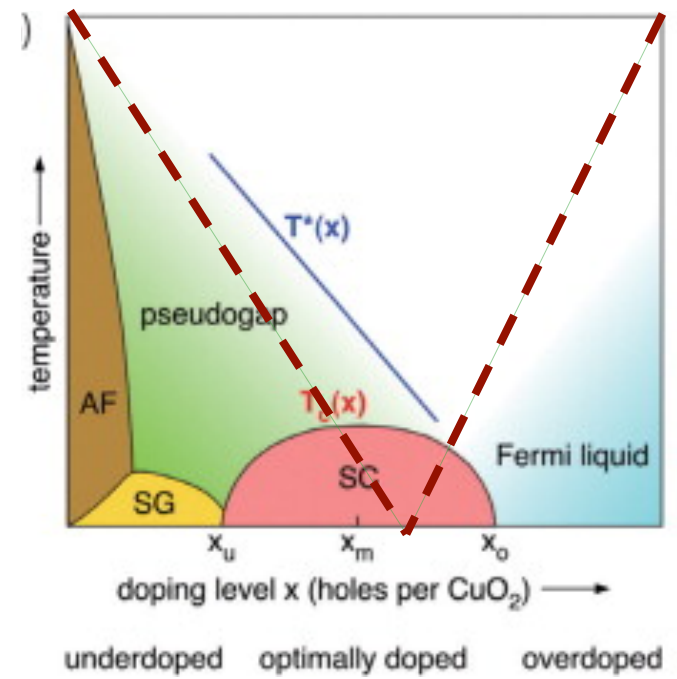


## Generic cuprate phase diagram

Bussman-Holder and Muller  
2008 review

# Cuprate Superconductivity

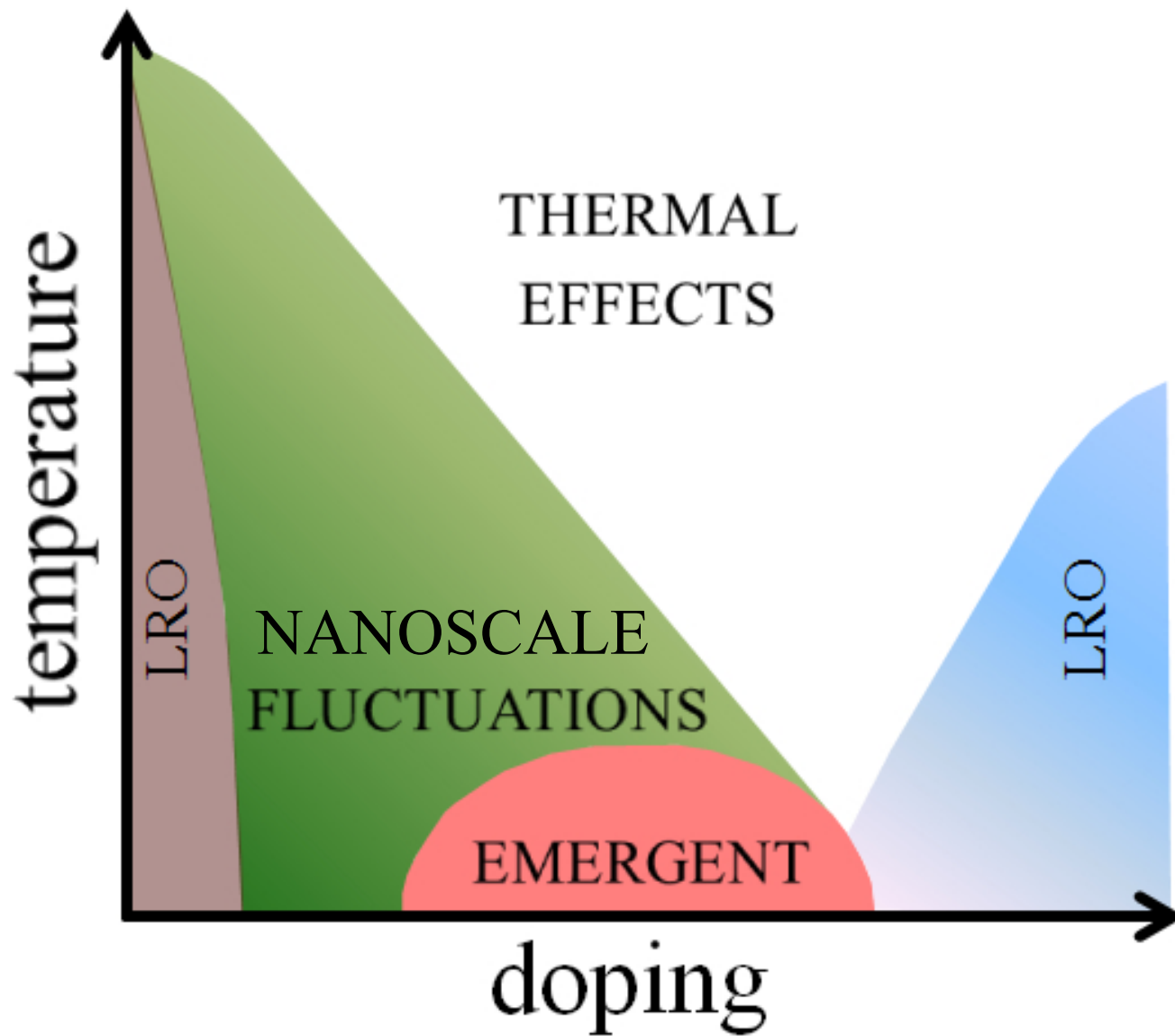
- Basic picture of doped, half-filled Mott insulator
- Antiferromagnetism and spin fluctuations
- Strange metal
- Spin gap  $\rightarrow$  pseudogap
- Is it an X or a V?
- Nanoscale electronic inhomogeneities (stripes)
  - Really hard to see! Took 8 years to find them in LBCO
  - Took 10 more years to find them in BSSCO (controversial STM results)
  - Took 7 more years to find them in YBCO (less controversial resonant scattering)



## Generic cuprate phase diagram

Bussman-Holder and Muller  
2008 review





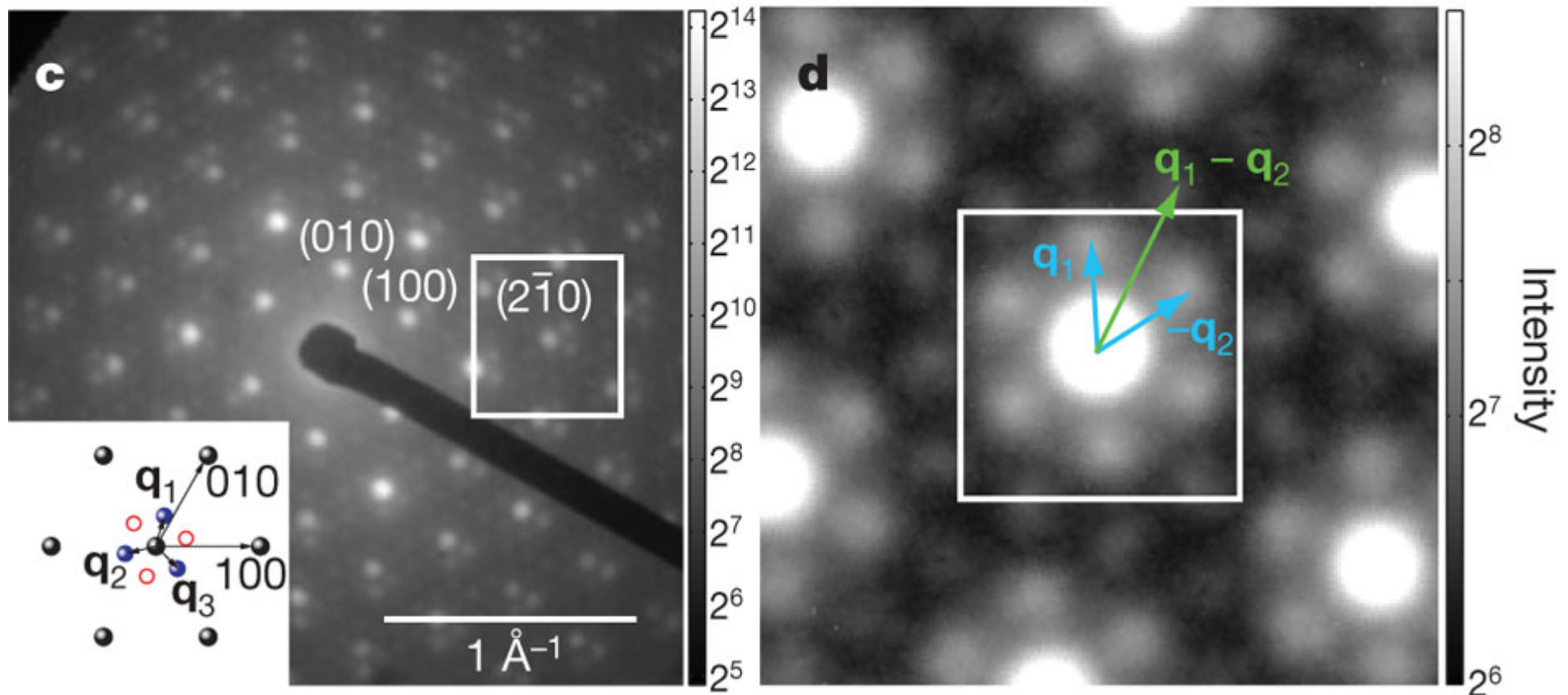
An even more generic phase diagram?????

# Materials with colossal emergent properties close to a metal-insulator transition

1. How can we detect and measure *LOCAL* broken symmetry electronic states in materials?
2. How ubiquitous are they?
3. How important are they?

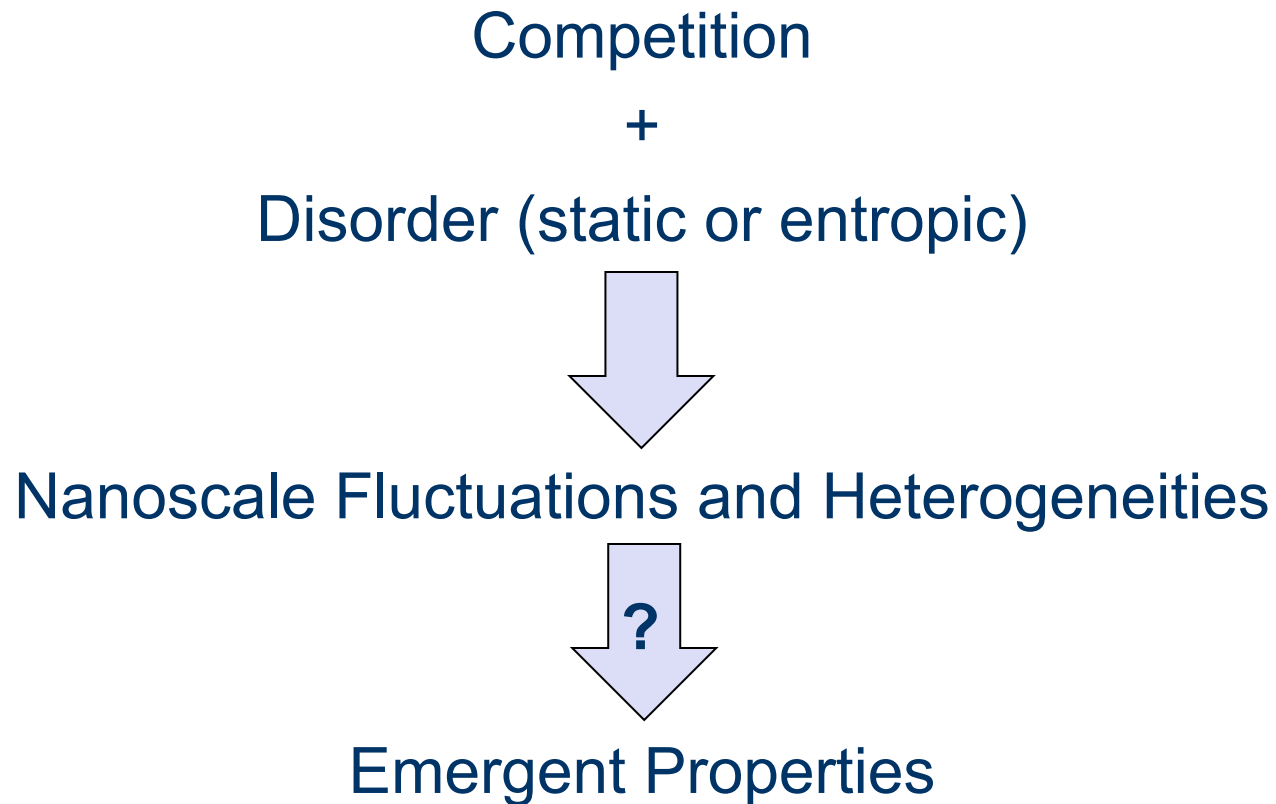
## How do we detect a global broken symmetry state?

- New Bragg Peaks! E.g., CDW in 1T-TaS<sub>2</sub>
- Eichberger et al Nature **468**, 799–802 (2010) doi:10.1038/nature09539
- Accompanying changes in transport/susceptibility/spectroscopy



# What happens if the broken symmetry state is only *short-range ordered*?

- Signal from the broken symmetry state is becomes diffuse
- Need a probe of local correlations
- Need to look for signals in measures of disorder in the material



# How can local broken symmetries hide?





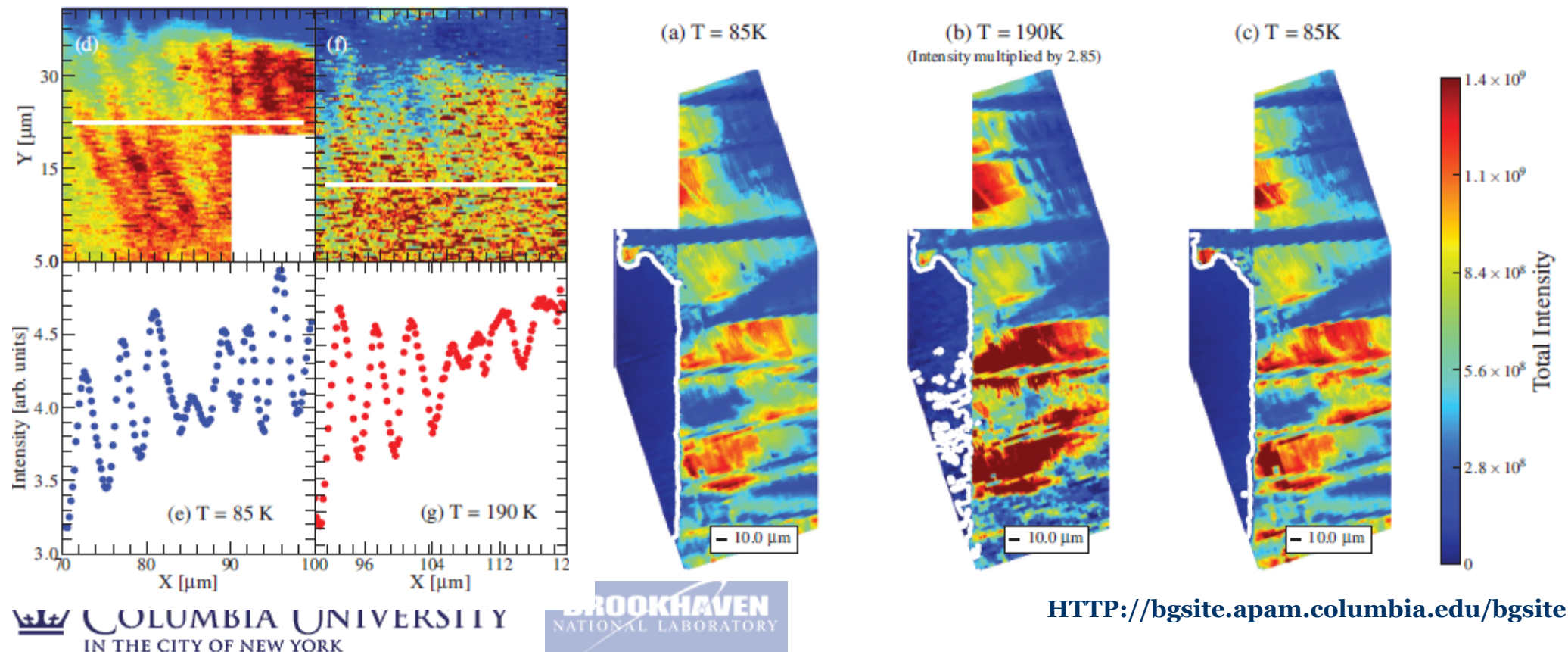
# How are hidden local broken symmetries revealed?

## #1: direct imaging



# Scientific Highlights: Nanoscale AF texture in manganites

- First spatially resolved images of nanoscale antiferromagnet texture using resonant soft x-ray scattering (PRB 2013)
- Applied to layered manganite materials



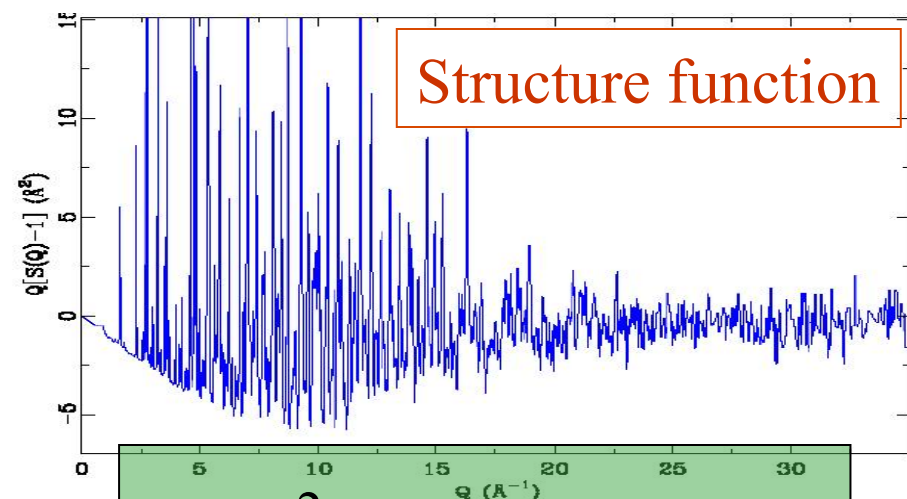
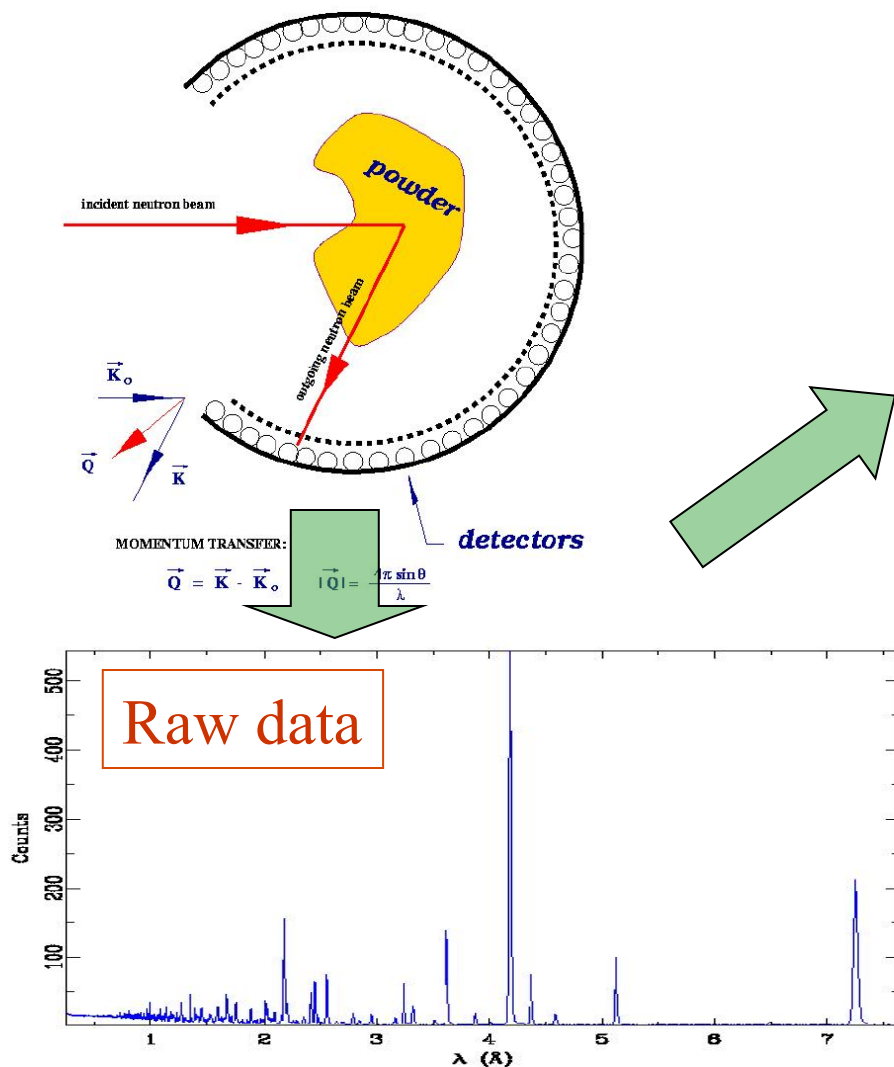


# How are hidden local broken symmetries revealed?

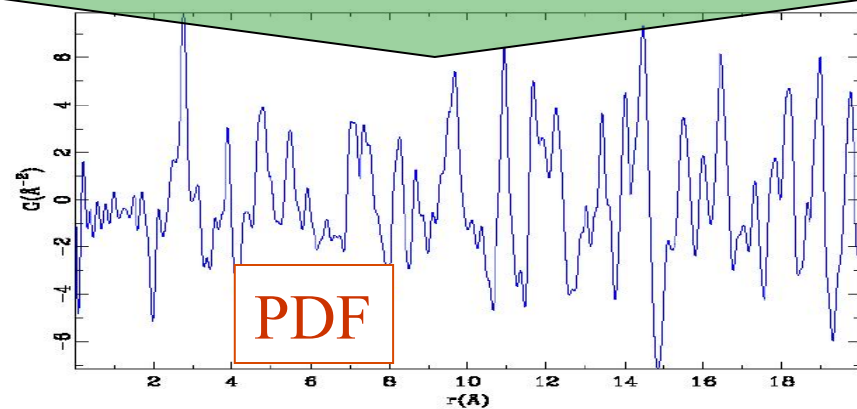
## #2 Studies of local structure/correlations



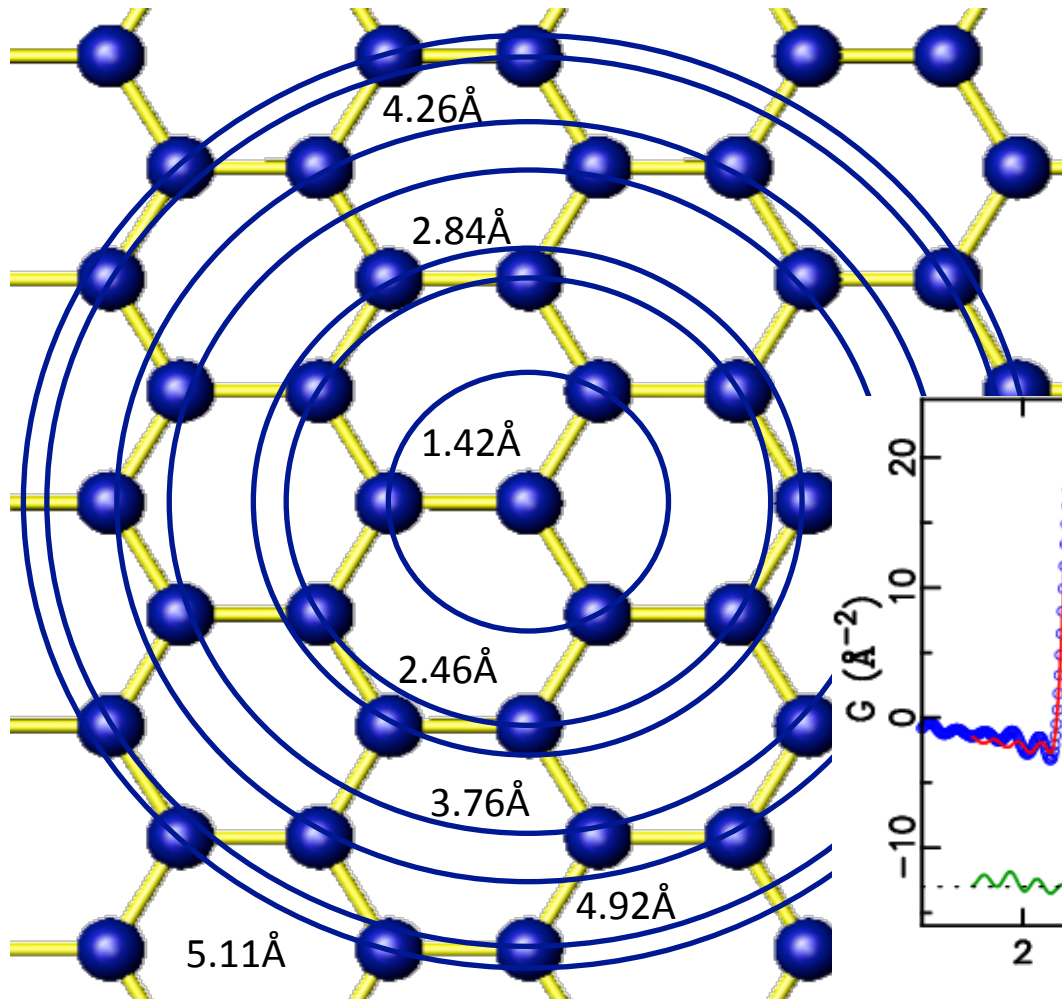
# Obtaining the PDF



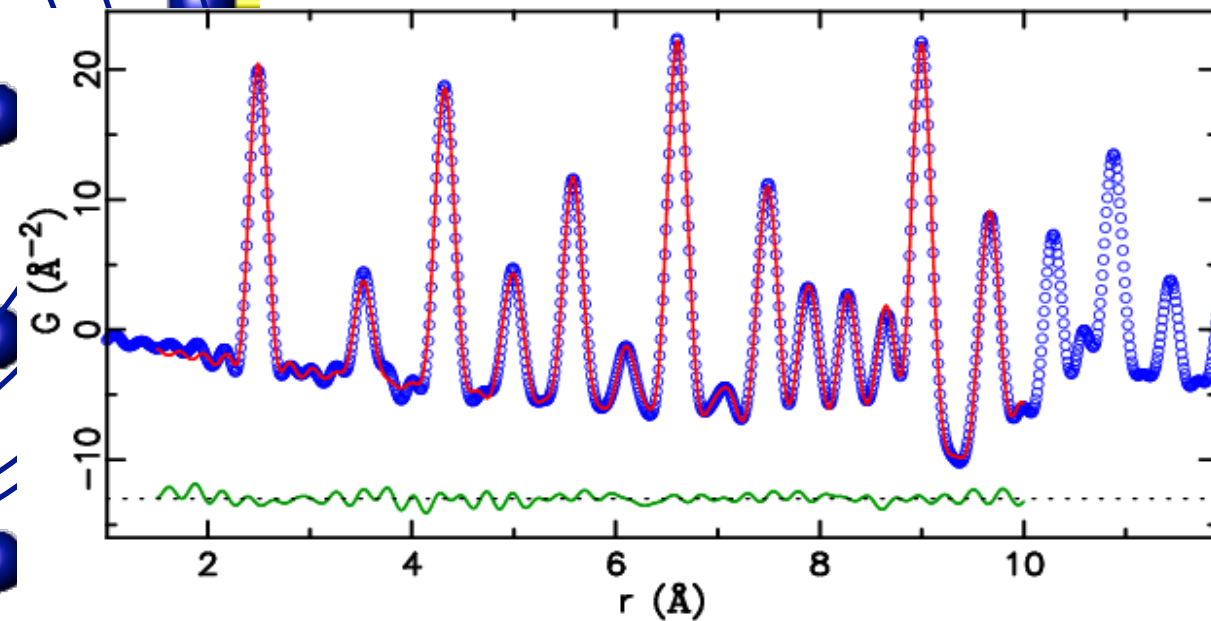
$$G(r) = \frac{2}{\pi} \int_0^\infty Q[S(Q)-1] \sin Qr dQ$$



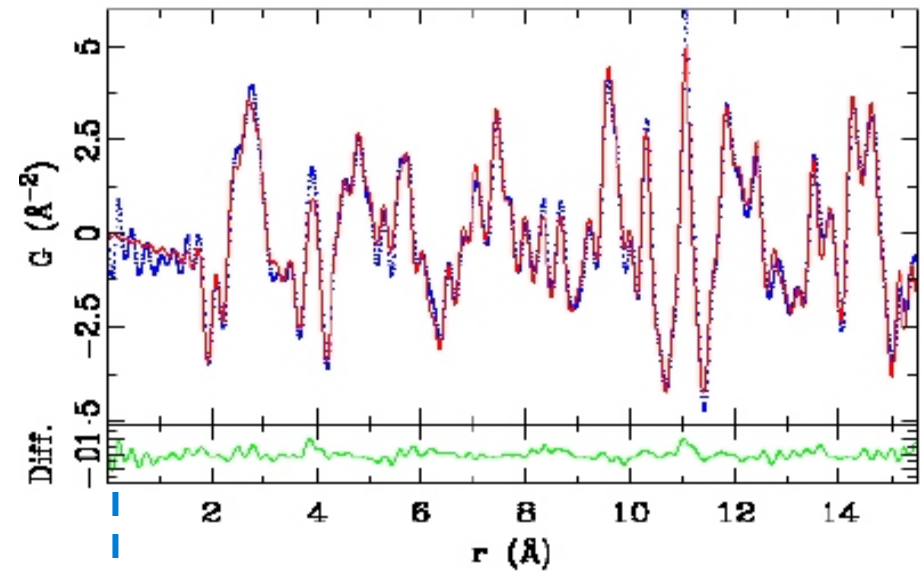
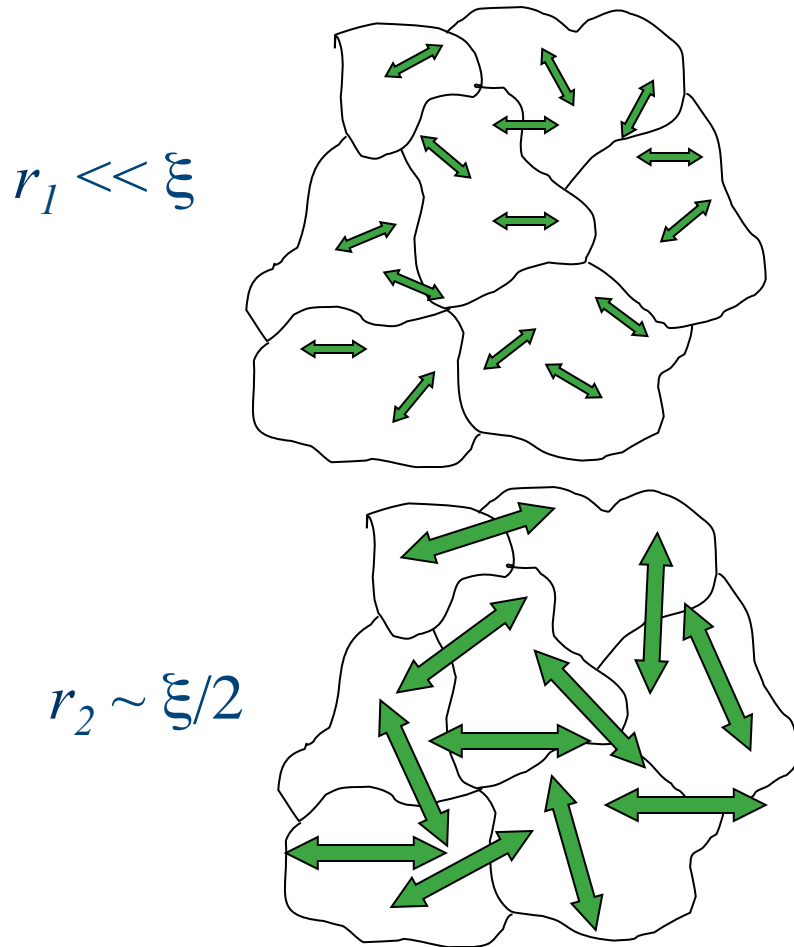
## Nanostructure refinement



Pair distribution function (PDF) gives the probability of finding an atom at a distance “r” from a given atom.



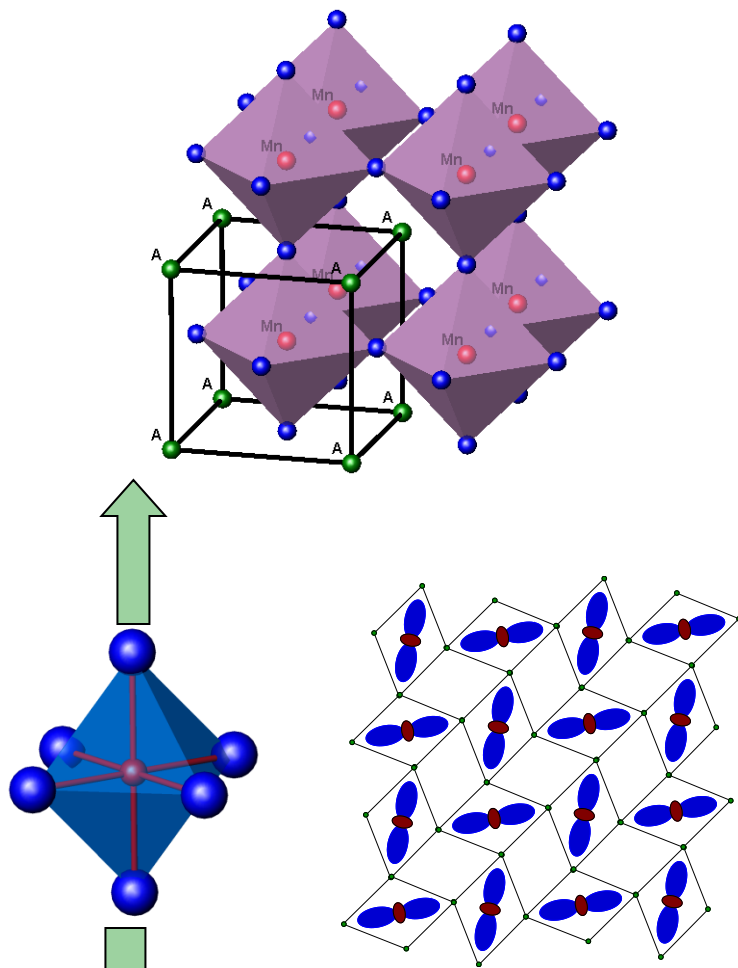
# Observing Nanodomains in the PDF



Intra-domain structure

Inter-domain structure

# How is PDF sensitive to MI transition?



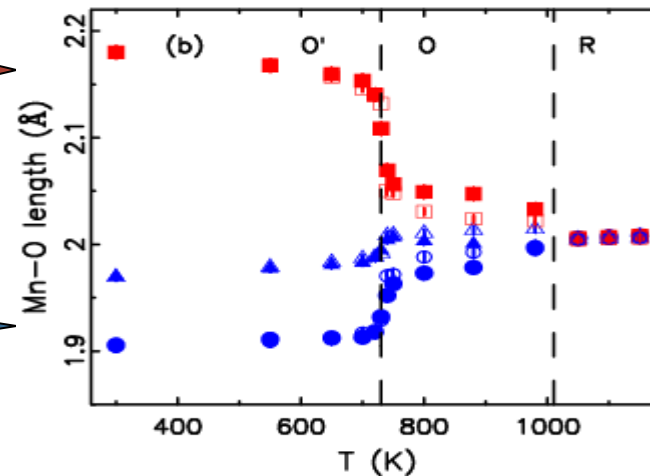
- **Mn<sup>3+</sup>:** JT distorted octahedra:
  - 4 x ~1.94 Å Mn-O bonds
  - 2 x 2.17 Å Mn-O bonds
  - $R_{\text{long}} - R_{\text{short}} = 0.23\text{Å}$
  - 8 long and 4 short O-O distances centered around 2.75Å
- **Mn<sup>4+</sup>:** undistorted octahedra:
  - 6 x ~1.94 Å Mn-O bonds
- **Delocalized holes:** undistorted octahedra:
  - 6 x ~1.94 Å Mn-O bonds

# T-dependence of Mn-O bond distribution

Long-bonds



Short-bonds



Average structure

- Mn-O bond lengths are invariant with temperature, right up into the R-phase

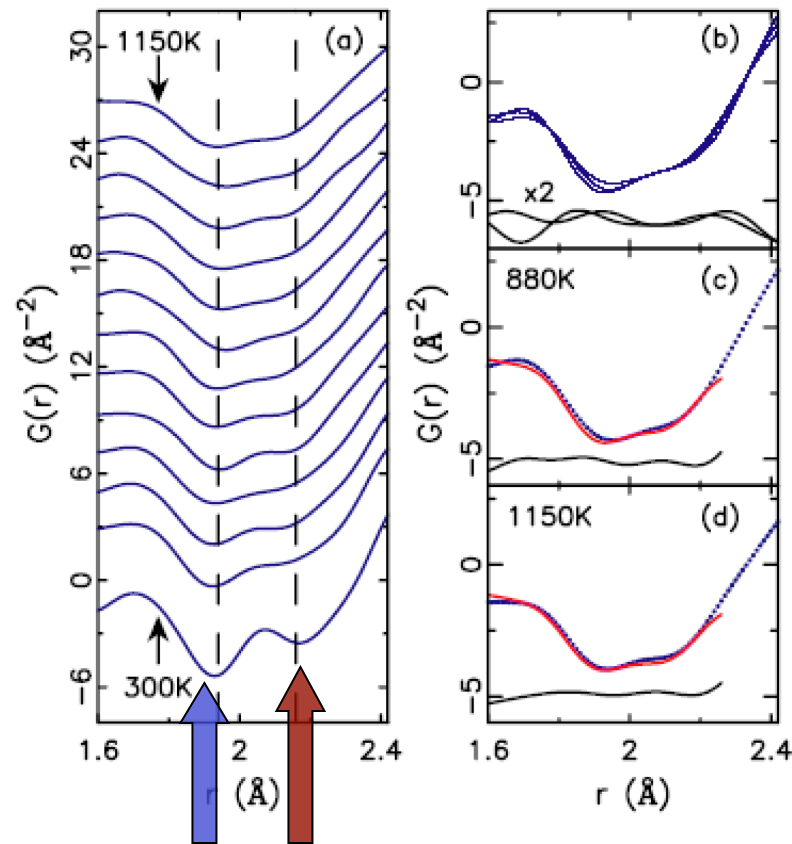
**JT distortions persist locally in the pseudocubic phase**

Xiangyun Qiu, Th. Proffen, J. F. Mitchell and S. J. L. Billinge, *Phys. Rev. Lett.* **94**, 177203 (2005).

Agrees with XAFS result: M. C. Sanchez et al., PRL (2003).



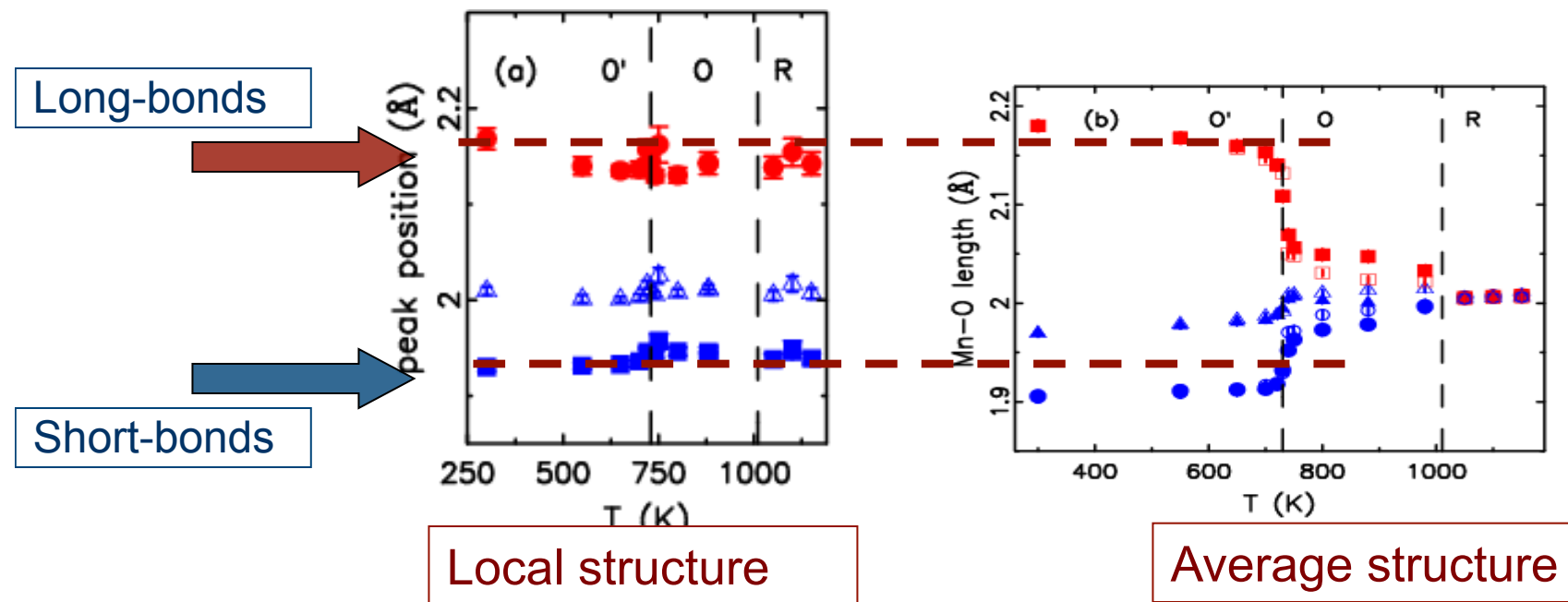
## Example: orbitally disordered phase in $\text{La}_{1-x}\text{Ca}_x\text{MnO}_3$



Local structure: JT distorted  
Average structure: no JT distortion  
=> Domains of local JT order

From Qui X., Proffen Th., Mitchell J. and Billinge S.J.L.  
PRL, 94, 177203 (2005).

# T-dependence of Mn-O bond distribution



- Mn-O bond lengths are invariant with temperature, right up into the R-phase

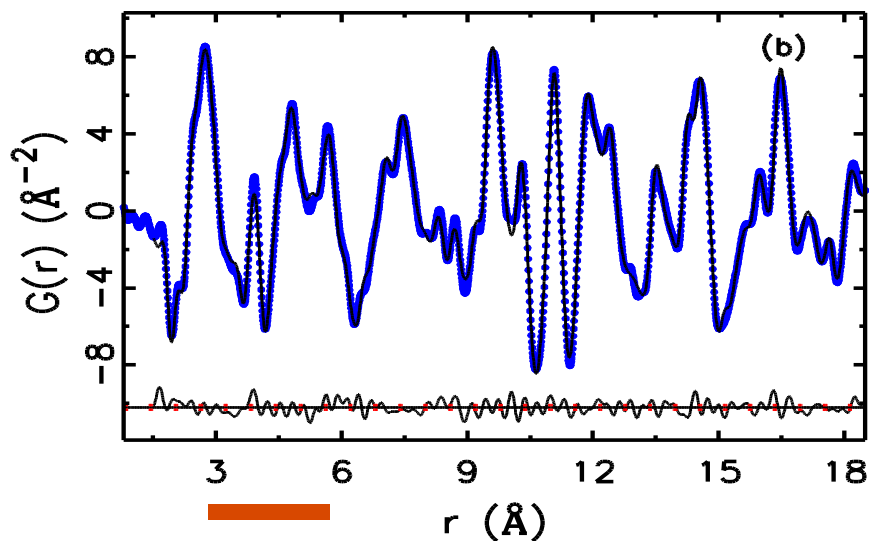
**JT distortions persist locally in the pseudocubic phase**

Xiangyun Qiu, Th. Proffen, J. F. Mitchell and S. J. L. Billinge, *Phys. Rev. Lett.* **94**, 177203 (2005).

Agrees with XAFS result: M. C. Sanchez et al., PRL (2003).



# Crossover from local to average structure

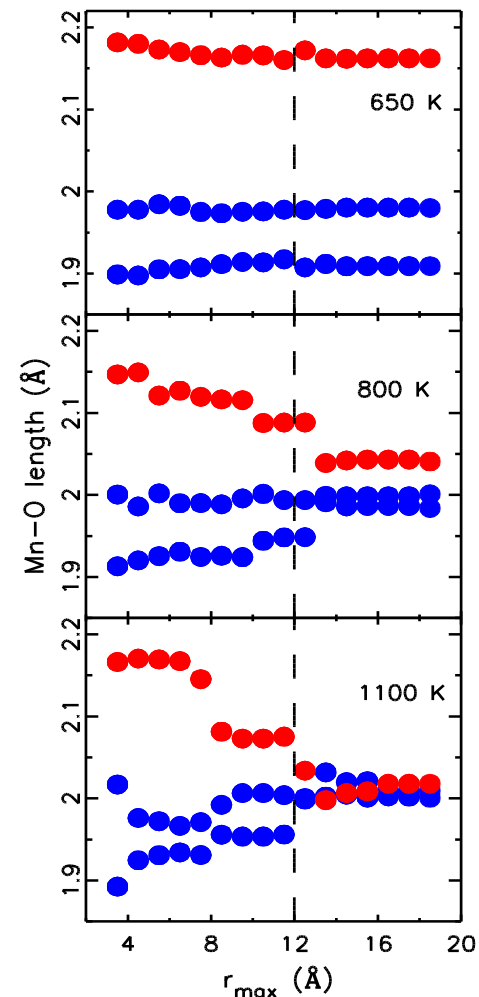


Local

Average

Intermediate???

- Varying range refinement
  - Fix  $r_{\min}$
  - Vary  $r_{\max}$
  - x axis is  $r_{\max}$

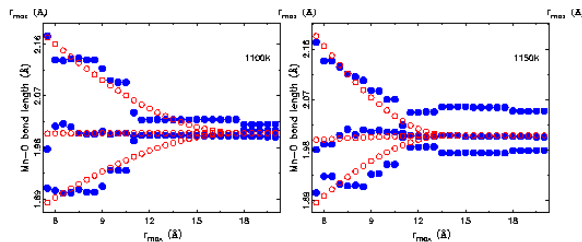
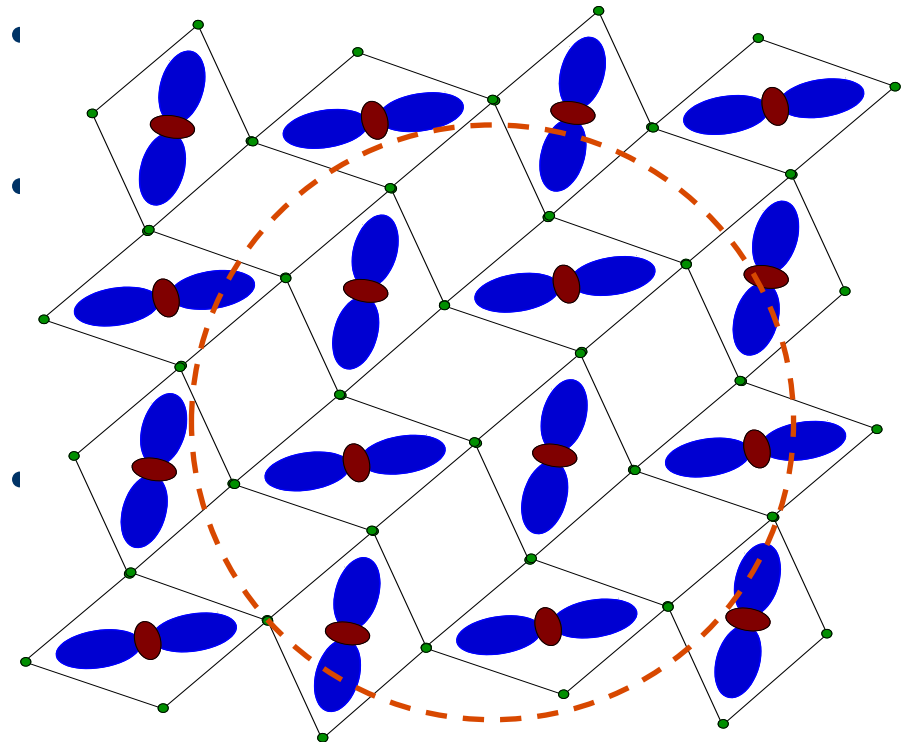
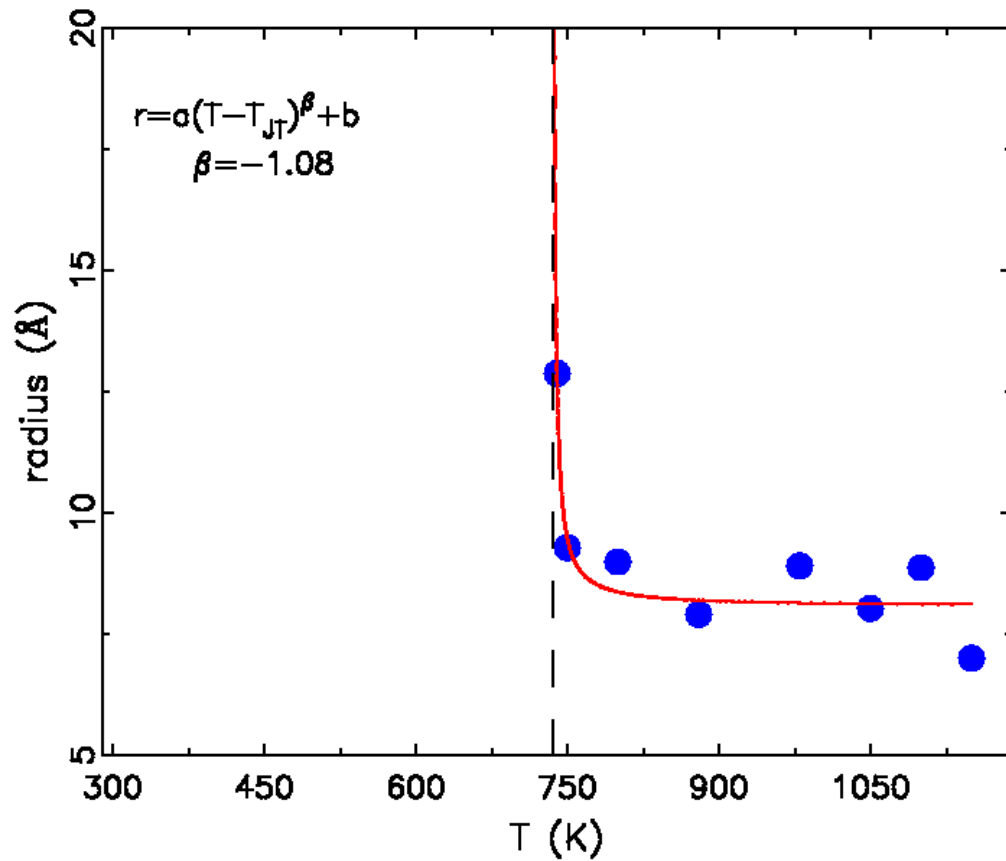


O

O'

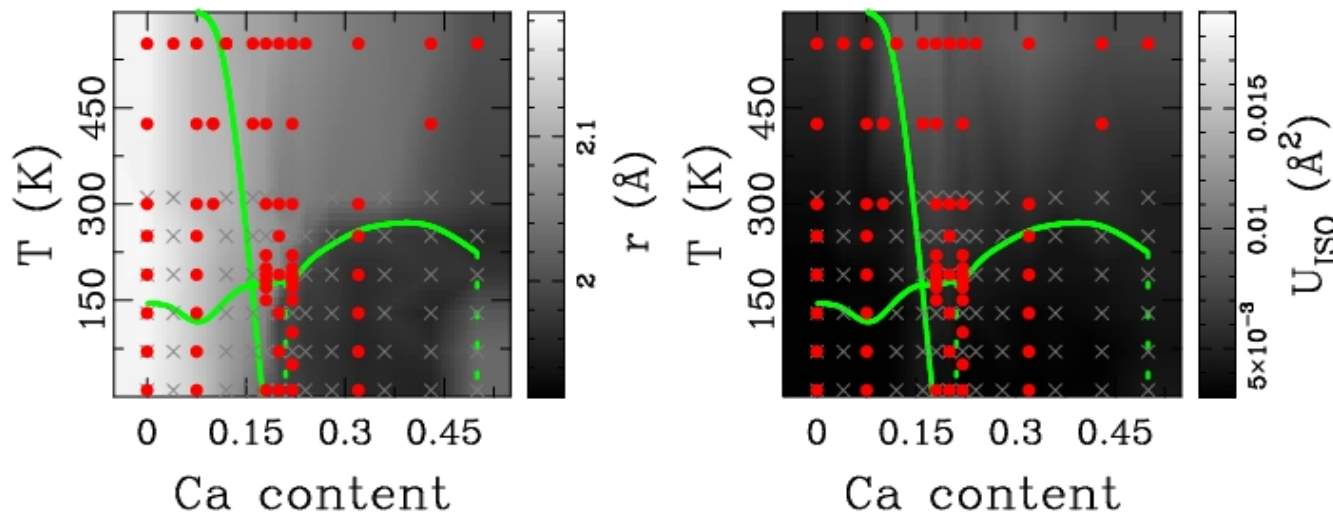
R

# R-dependent fits vs T



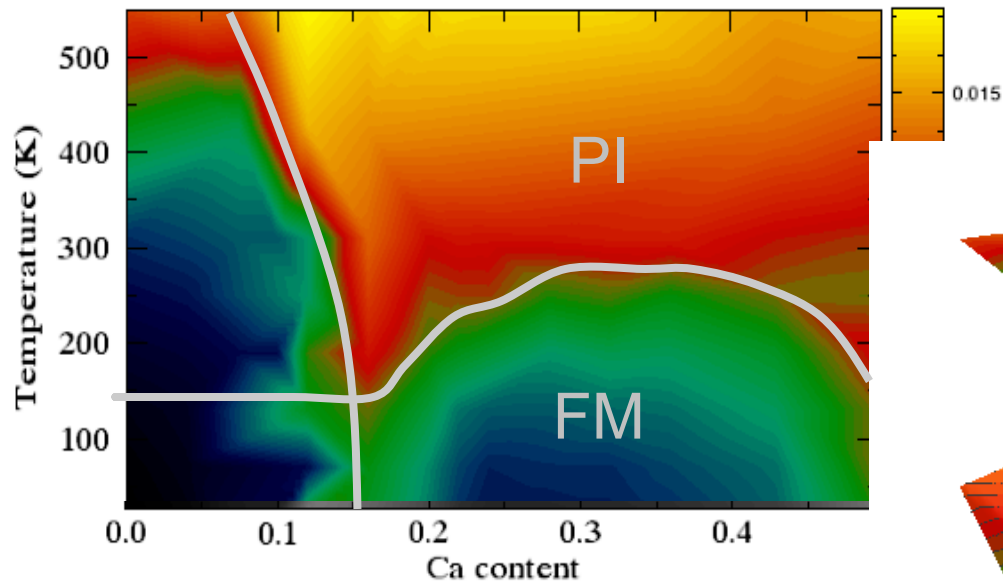
$$f(r, d) = \left[ 1 - \frac{3r}{2d} + \frac{1}{2} \left( \frac{r}{d} \right)^3 \right] \Theta(d - r)$$

# $\text{La}_{1-x}\text{Ca}_x\text{MnO}_3$ and colossal magnetoresistance

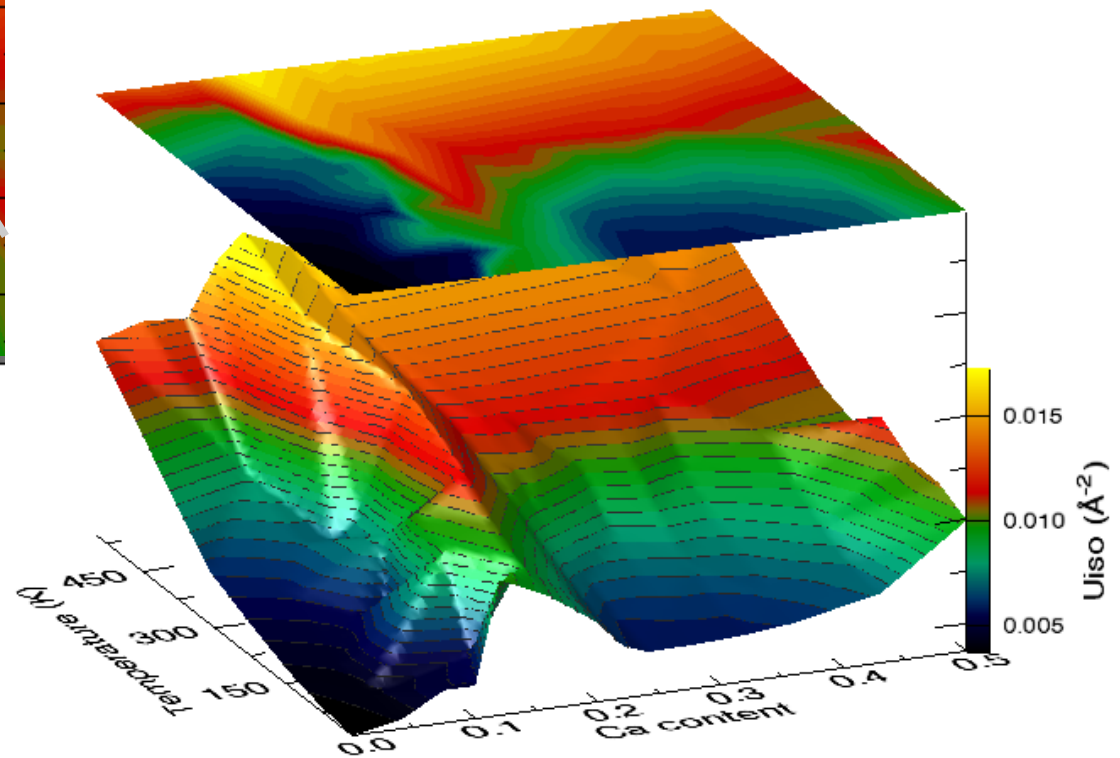


- Key point is that understand colossal magnetoresistance, we need many data points throughout the phase diagram.
- High throughput local structural studies!

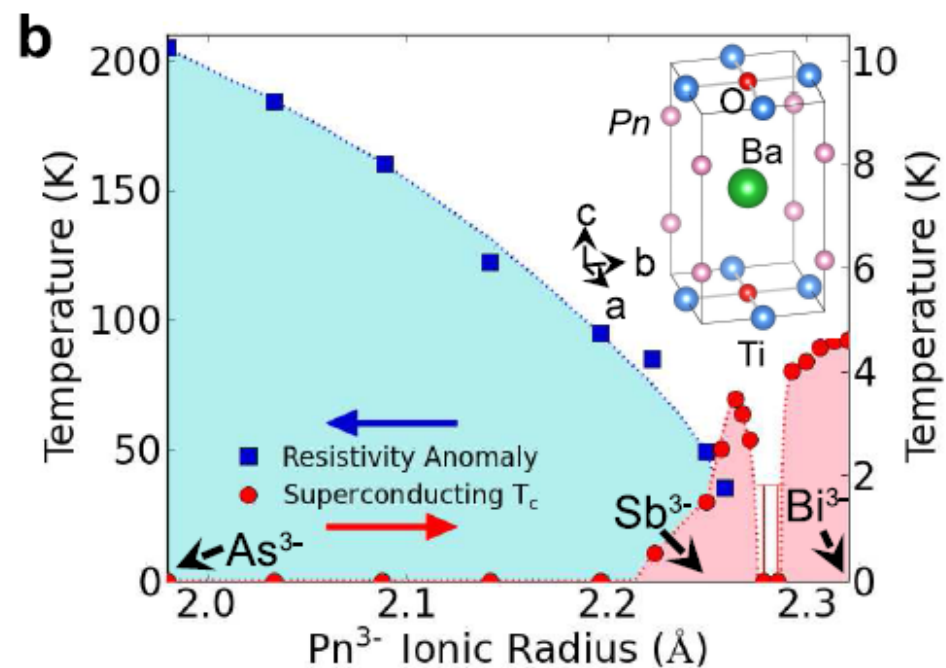
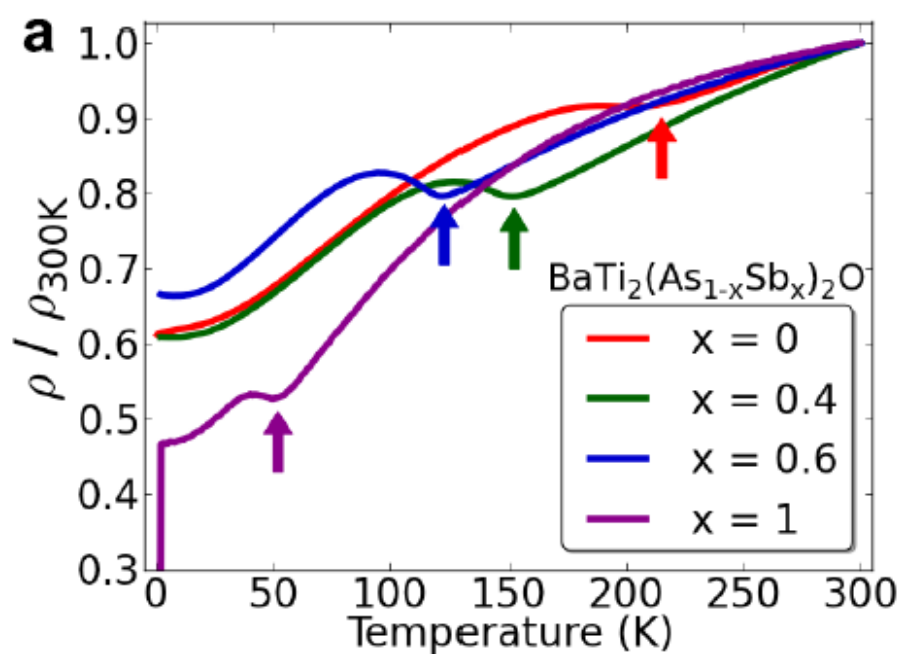
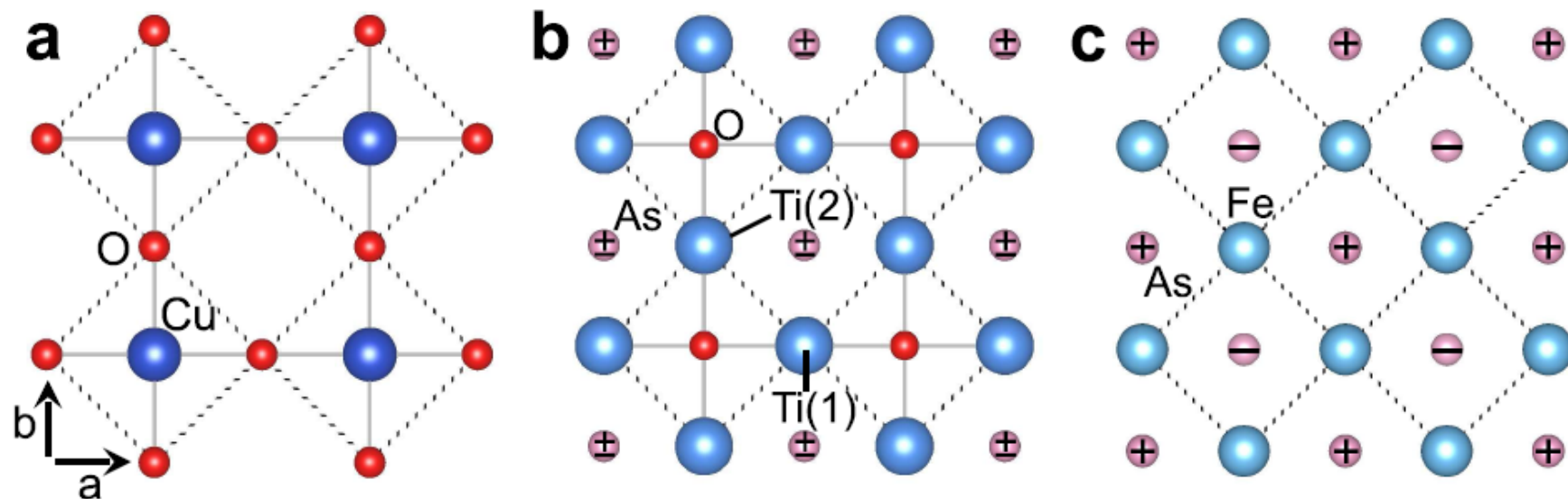
## Remarkable dataset: Phase diagram emerges from single refined parameters



- Phase diagram draws itself from the parameters
- Unexpected detail emerges and demands interpretation

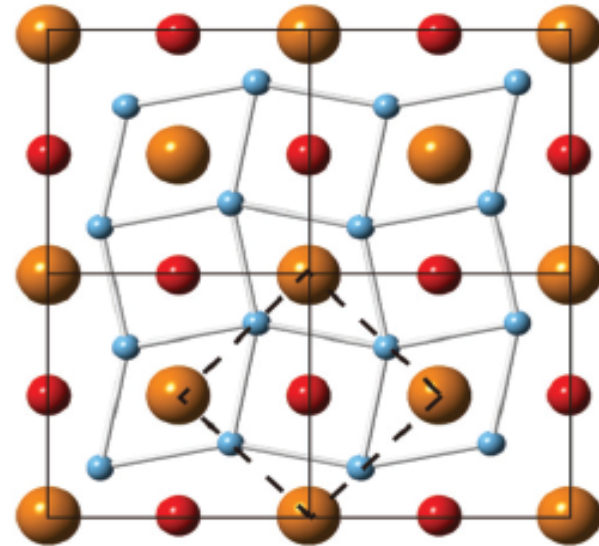


# BaTi<sub>2</sub>Pn<sub>2</sub>O (Pn = As, Sb, Bi)

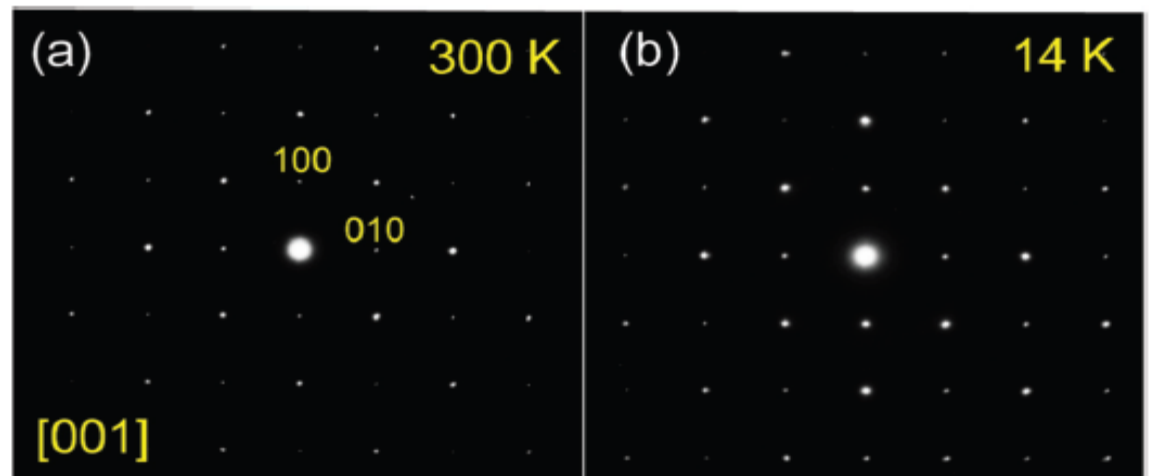


## Searching for the CDW

Proposed CDW distortion due to unstable phonon mode calculated from first principles (Subedi, PRB 87: 054506 (2013))



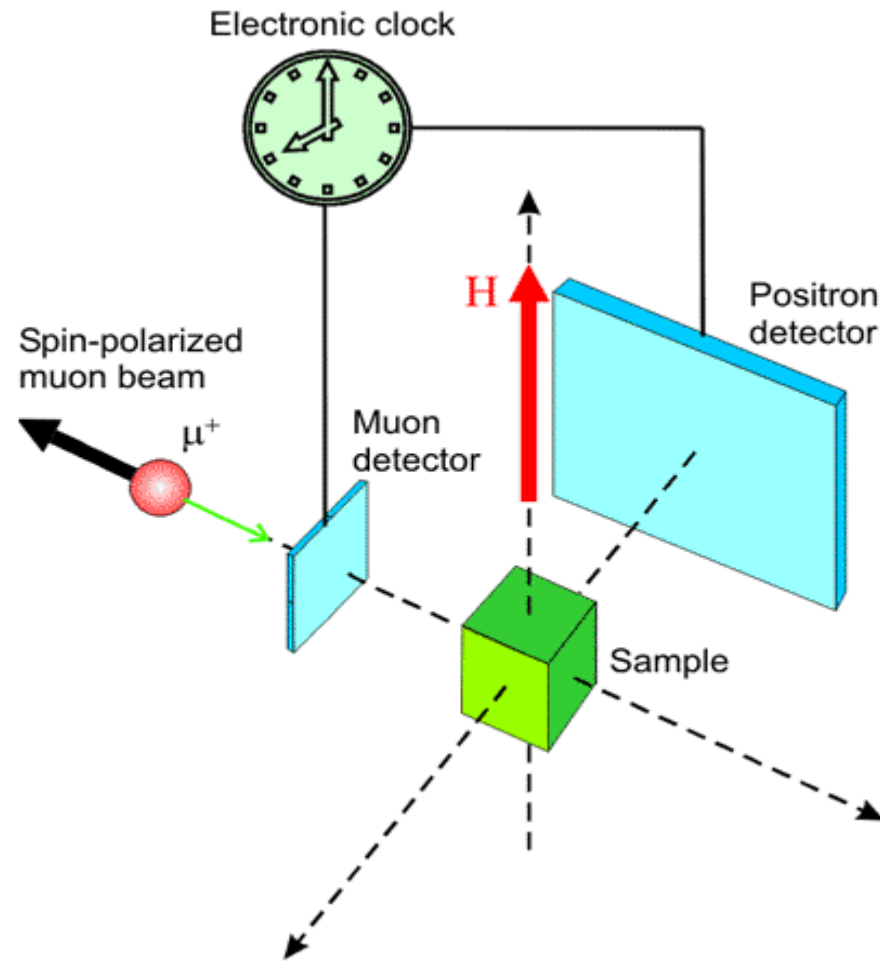
Observed electron diffraction patterns from  $\text{BaTi}_2\text{Sb}_2\text{O}$  and  $\text{BaTi}_2\text{As}_2\text{O}$ : no superlattice peaks!



BOOKHAVEN  
ONAL LABORATORY

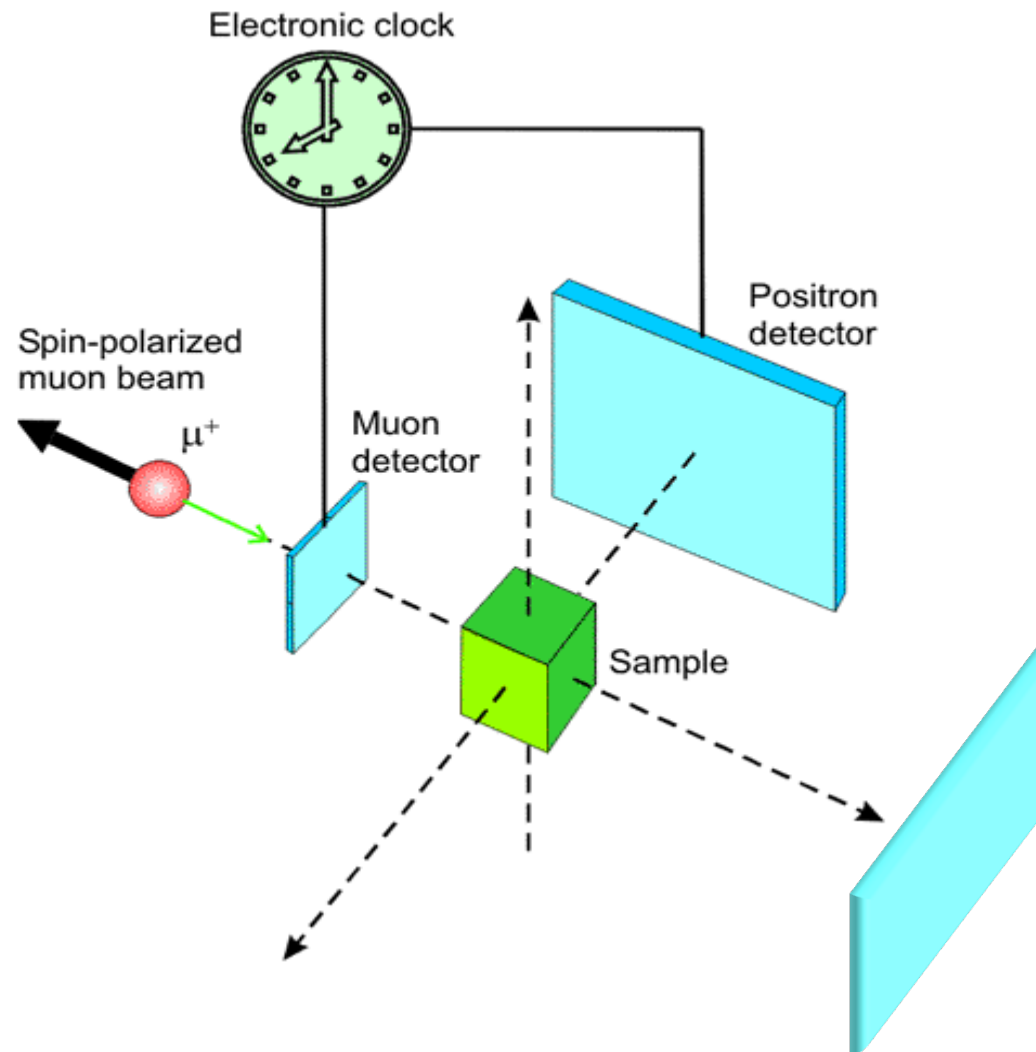
[HTTP://bgsite.apam.columbia.edu/bgsite](http://bgsite.apam.columbia.edu/bgsite)

# Probing magnetism with zero-field $\mu$ SR



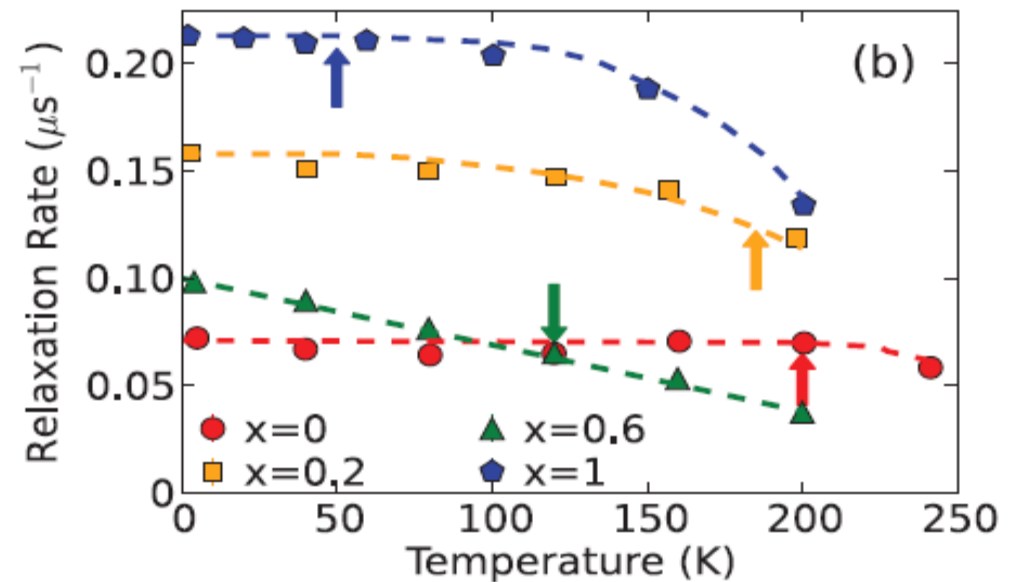
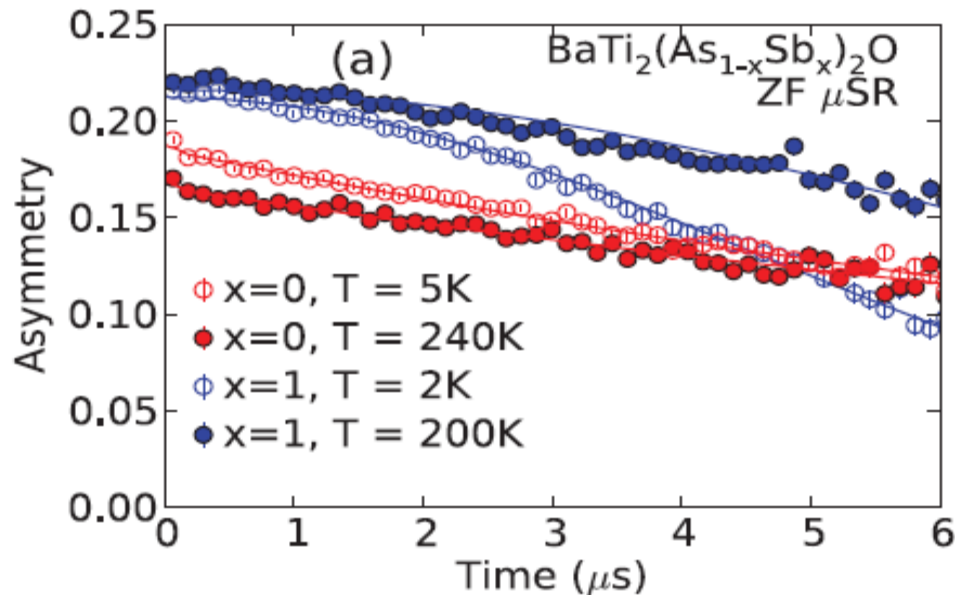


# Probing magnetism with zero-field $\mu$ SR





## Magnetic order probed by $\mu$ SR

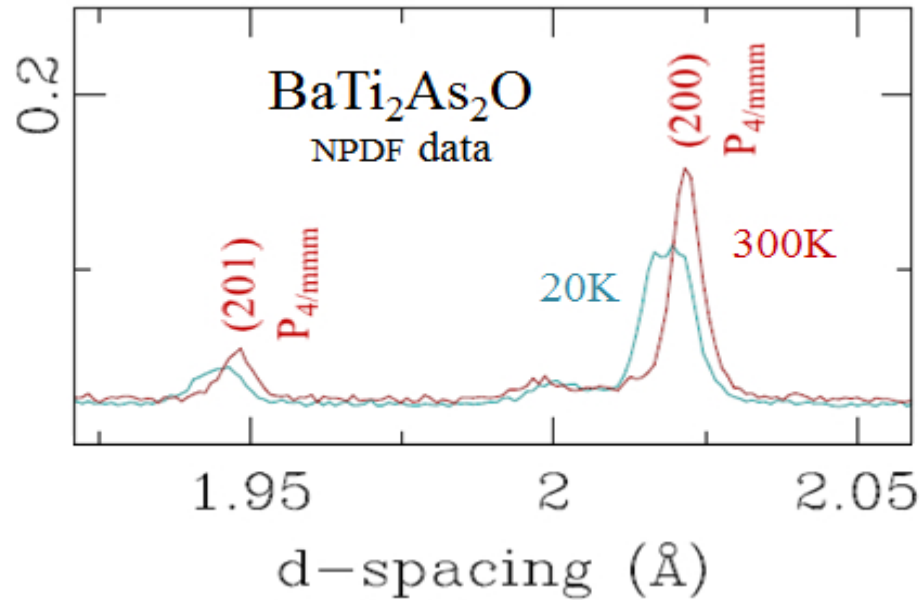


Y. Nozaki et al. PRB 88 214506 (2013)

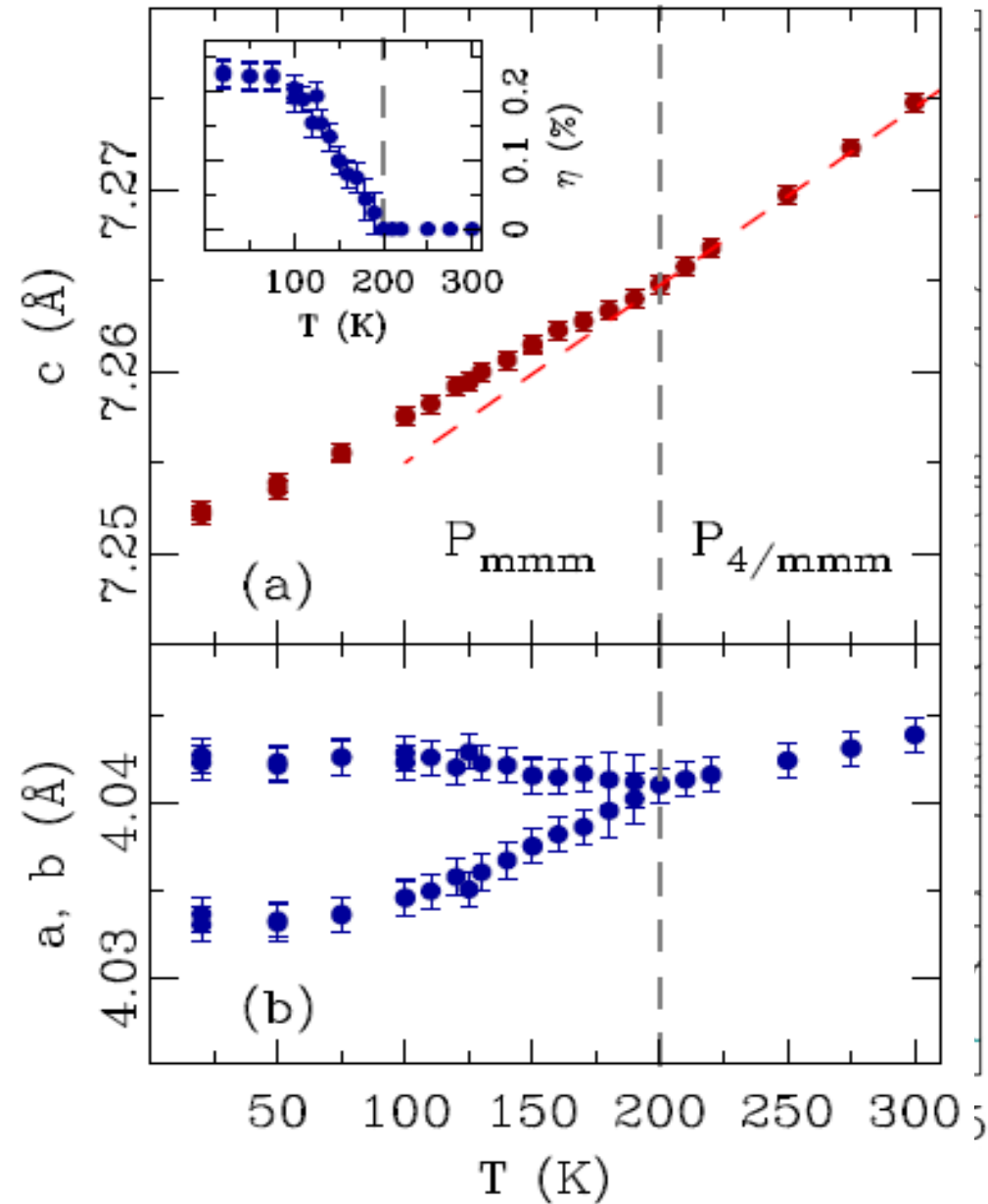
Lack of increase in relaxation rate at density-wave temperature indicates no magnetic order

Conclusion: No SDW in  $\text{BaTi}_2(\text{As,Sb})_2\text{O}$

Result is (probably) nematic charge order



Subtle but observable orthorhombic splitting of (200) and (201) peaks starting below approximately 200 K



## Case study #2: Dimers in $\text{CuIr}_2\text{S}_4$

Why is it interesting?

1. Frustration
2. Metal-insulator transitions
3. Carrier localization and broken symmetry states at low-temperature
4. Signal from broken symmetry state is large (colossal)!

# $\text{CuIr}_2\text{S}_4$ , Why is it interesting?

1. Frustration
2. Metal-insulator transitions
3. Carrier localization and broken symmetry states at low-temperature
4. Signal from broken symmetry state is large (colossal)!

# Spinel structure

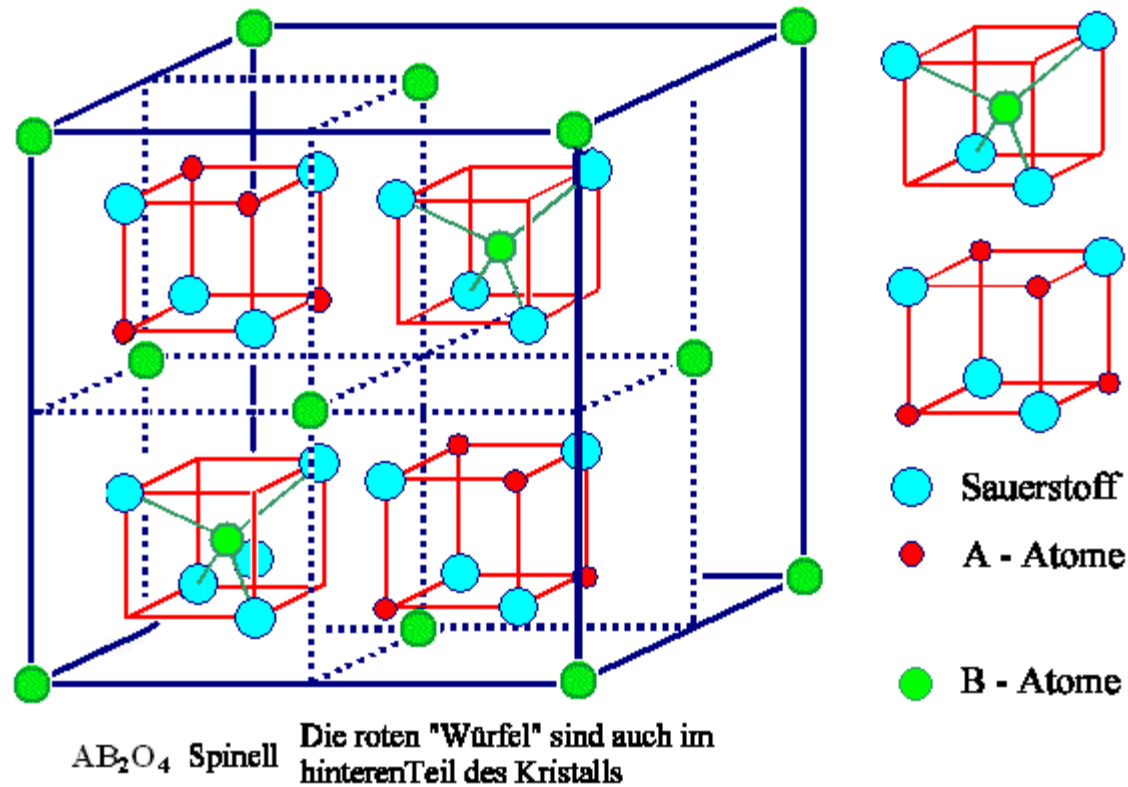
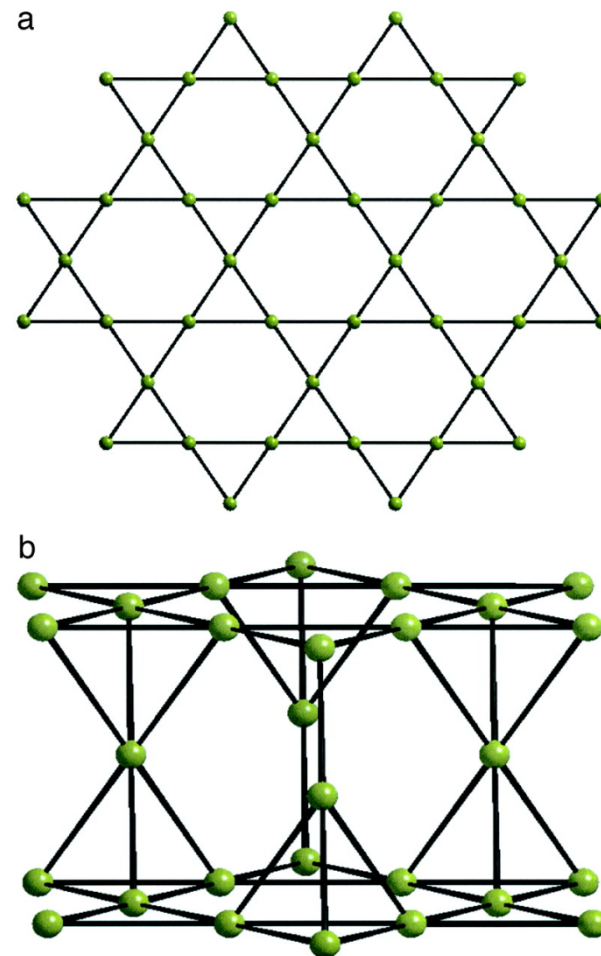
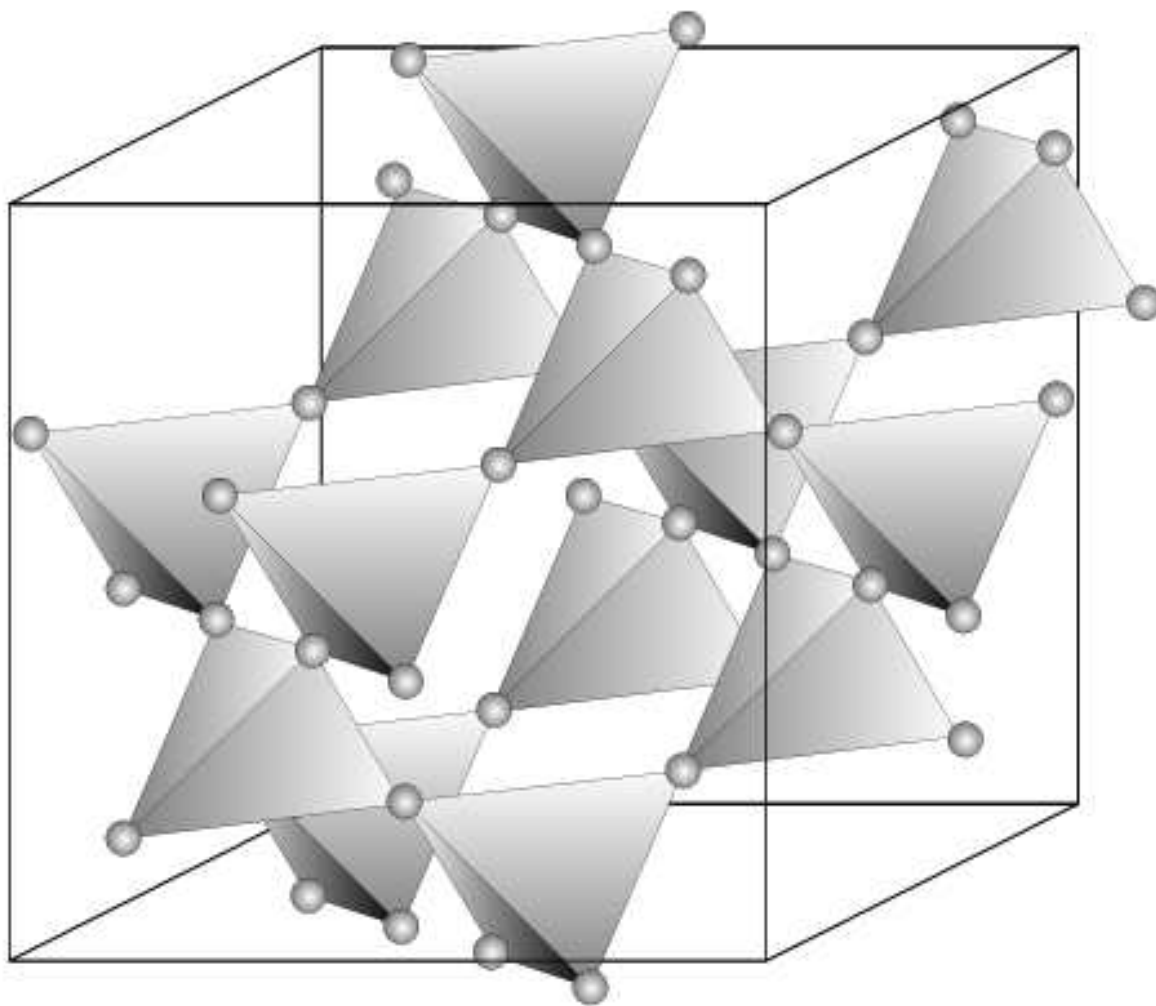
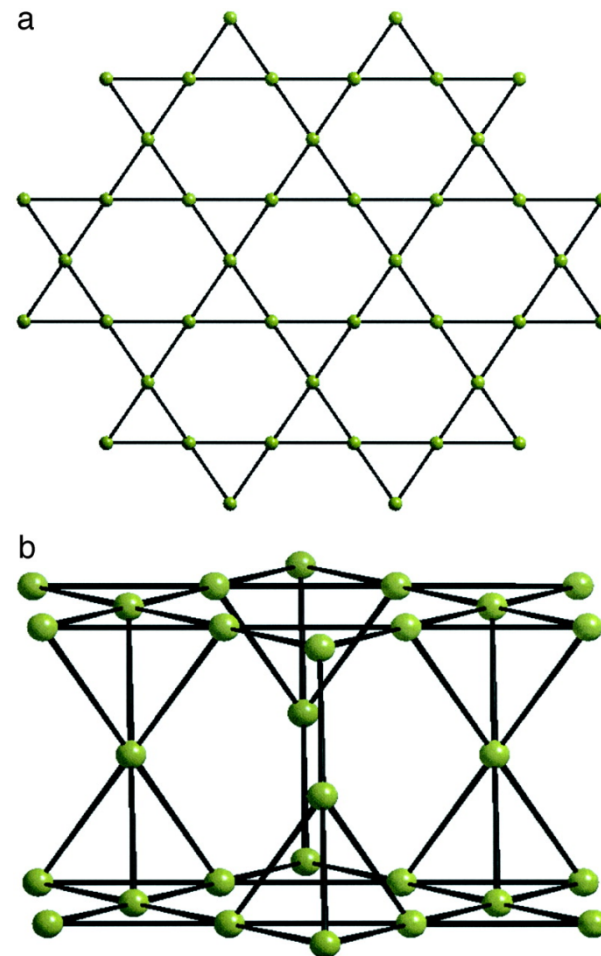
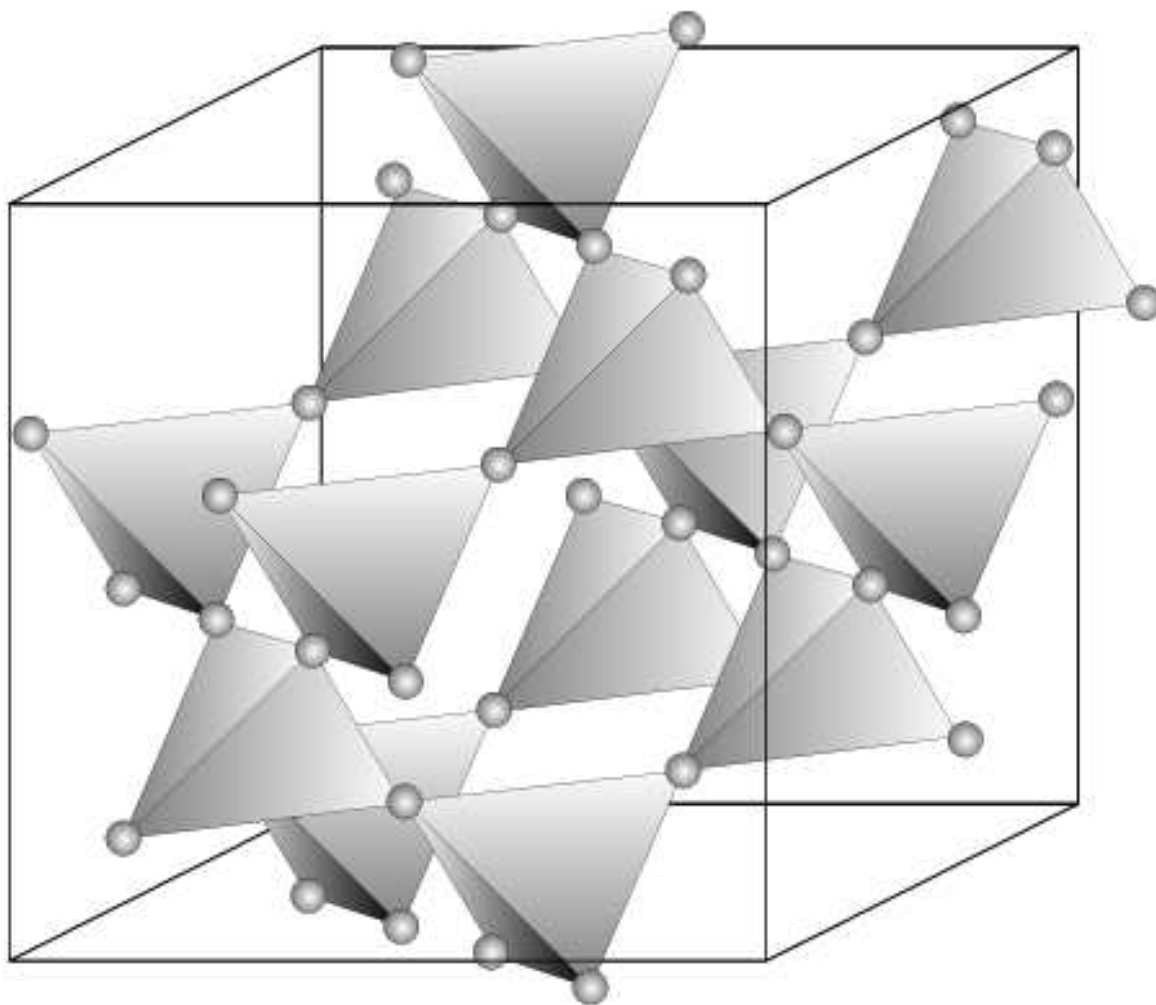


Image credit: [http://www.tf.uni-kiel.de/matwis/amat/def\\_en/kap\\_2/basics/b2\\_1\\_6.html](http://www.tf.uni-kiel.de/matwis/amat/def_en/kap_2/basics/b2_1_6.html)

But the B atoms form a pyrochlore lattice of corner shared tetrahedra => frustration

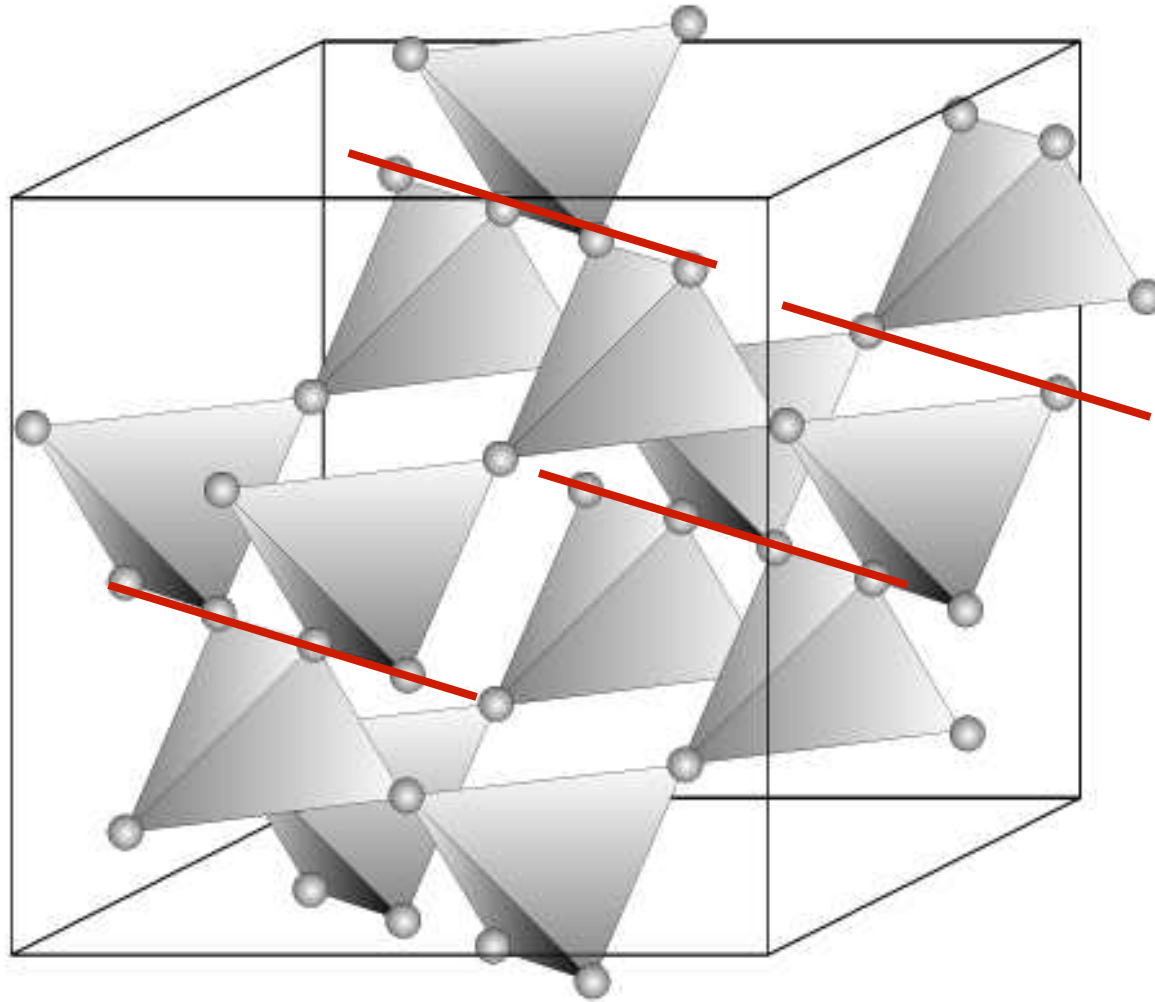


But the B atoms form a pyrochlore lattice of corner shared tetrahedra => ~~frustration~~

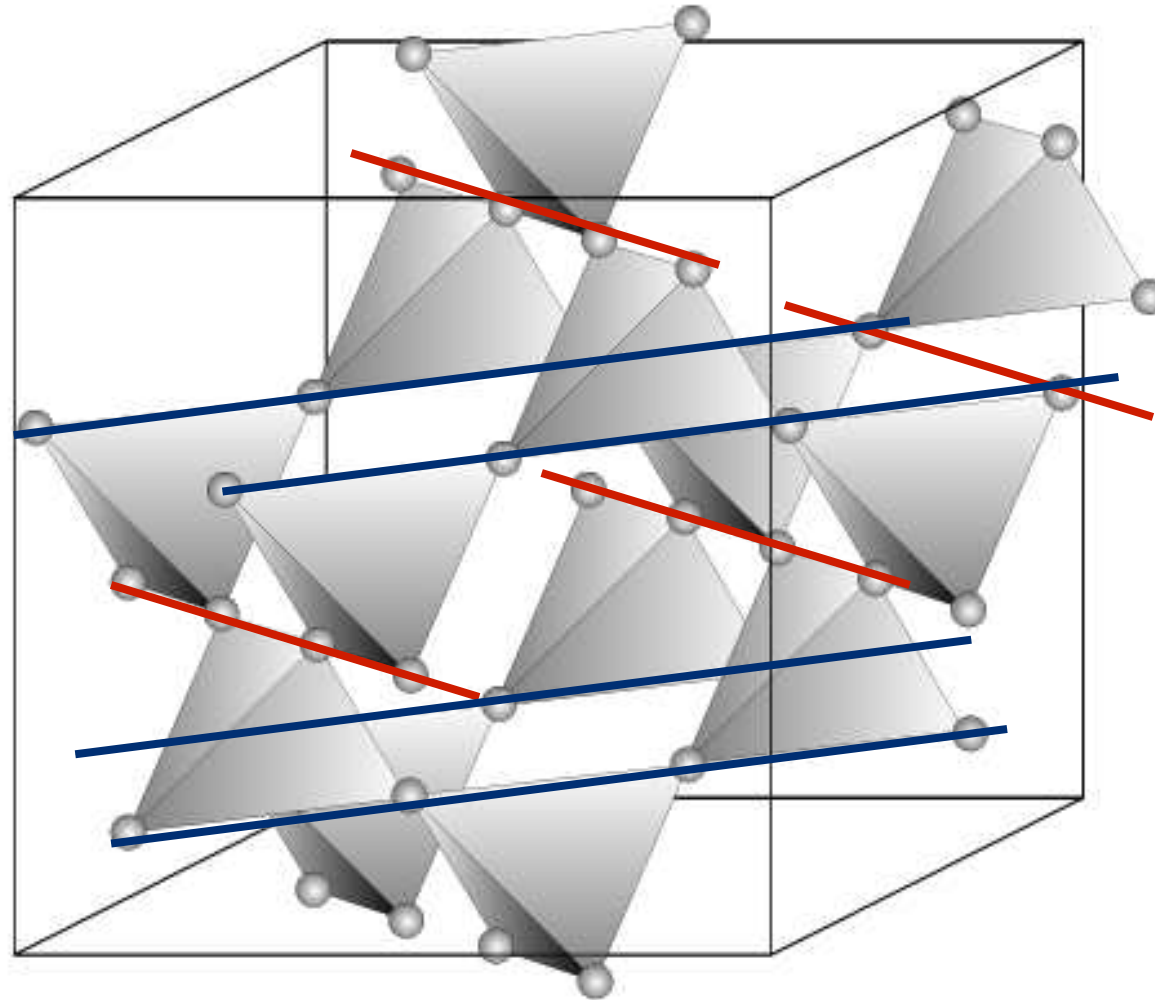




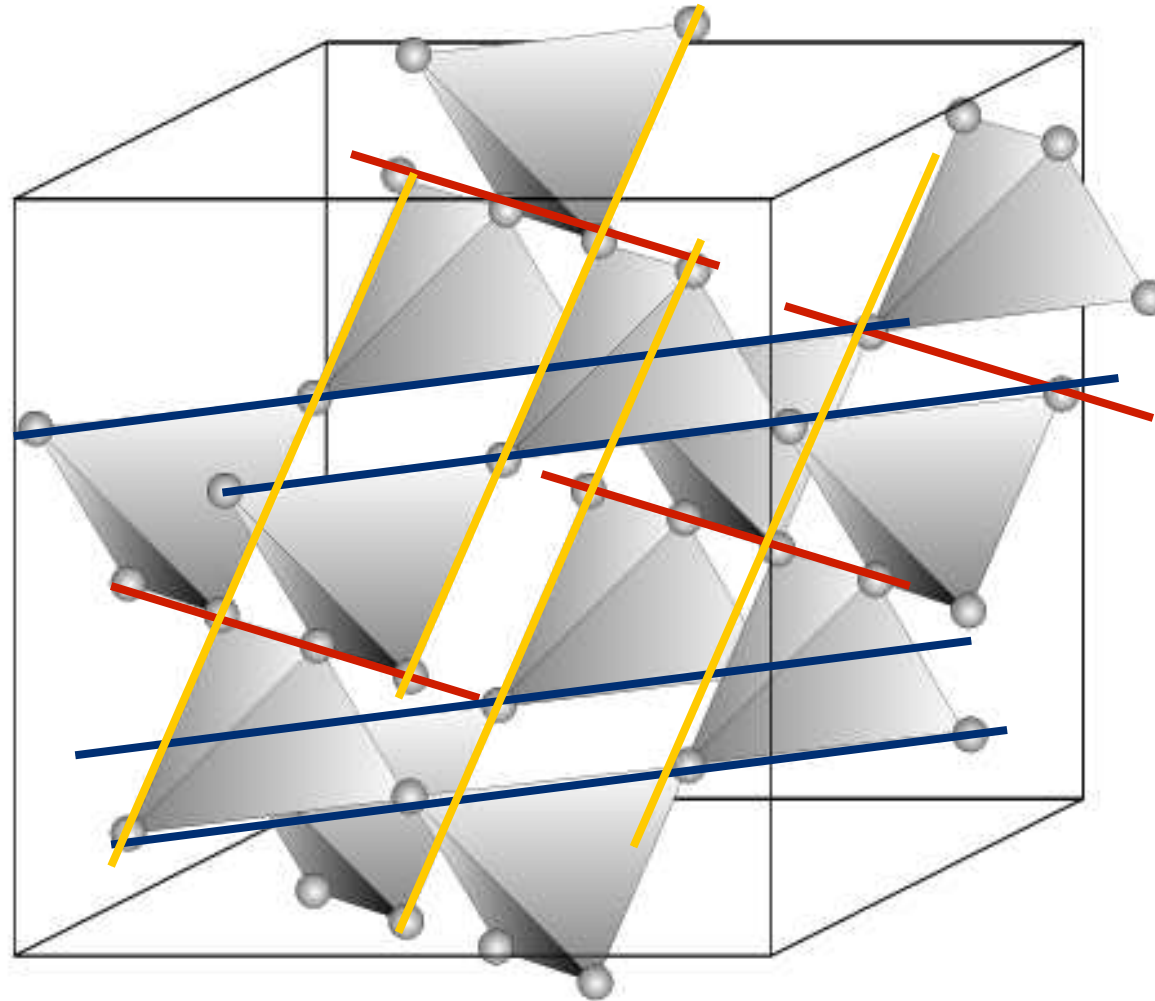
But the B atoms form a pyrochlore lattice of corner shared tetrahedra => low dimensional physics!

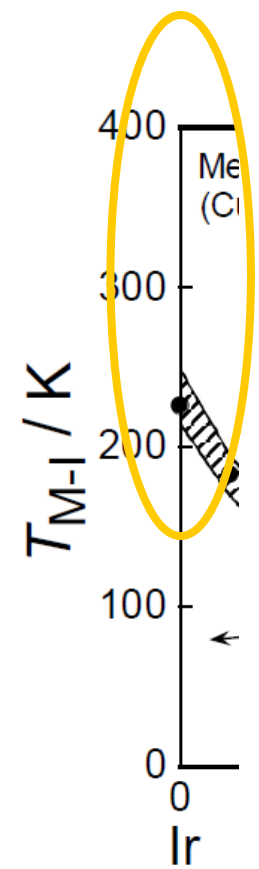
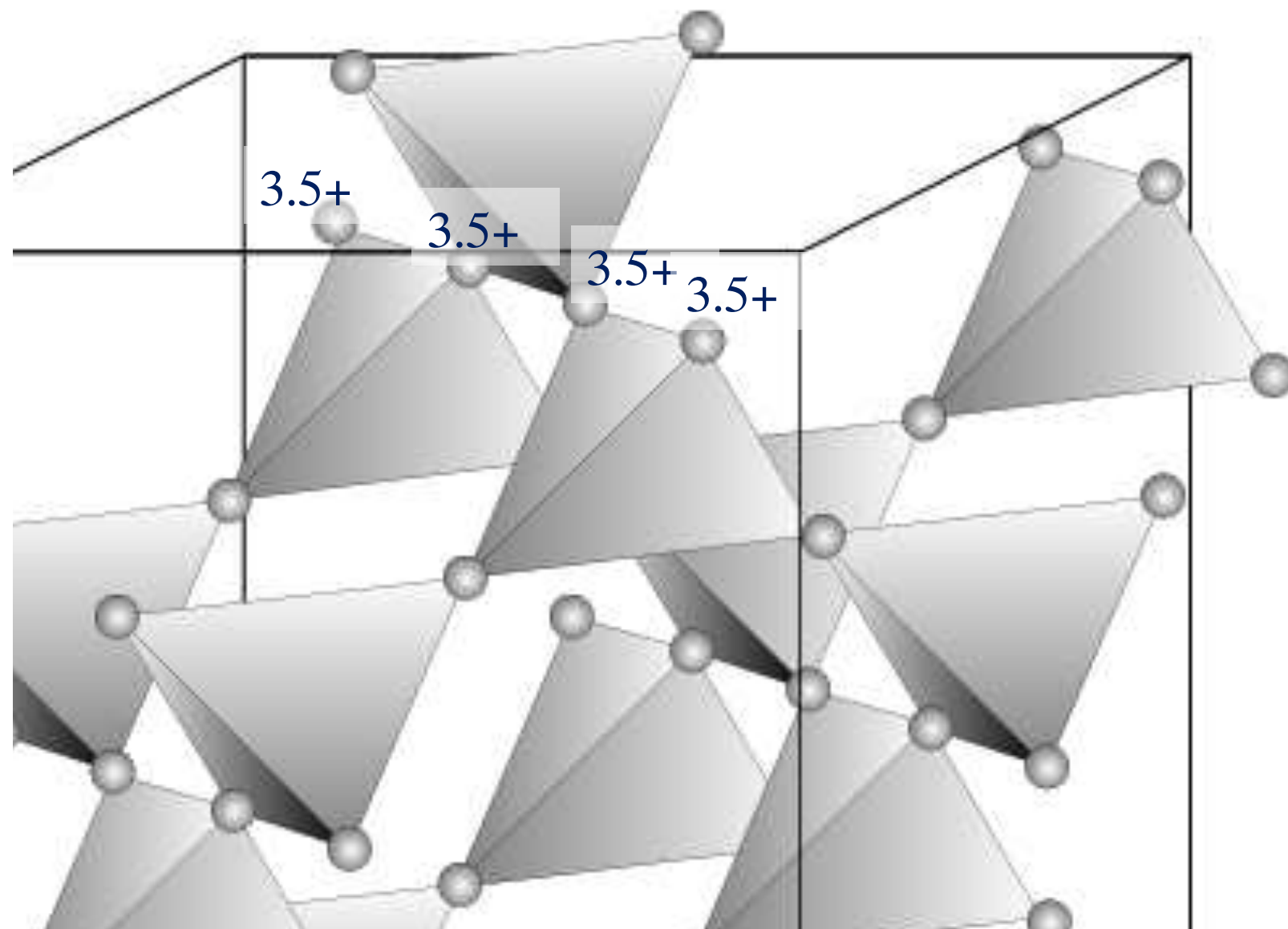


But the B atoms form a pyrochlore lattice of corner shared tetrahedra => low dimensional physics!



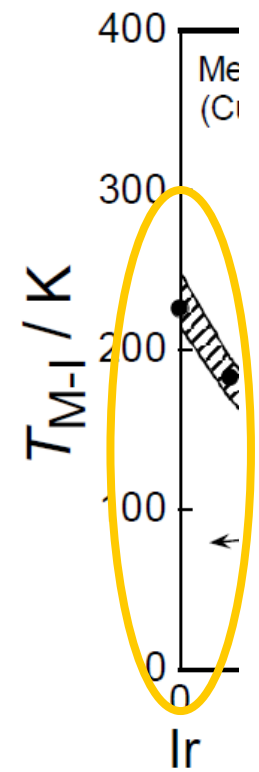
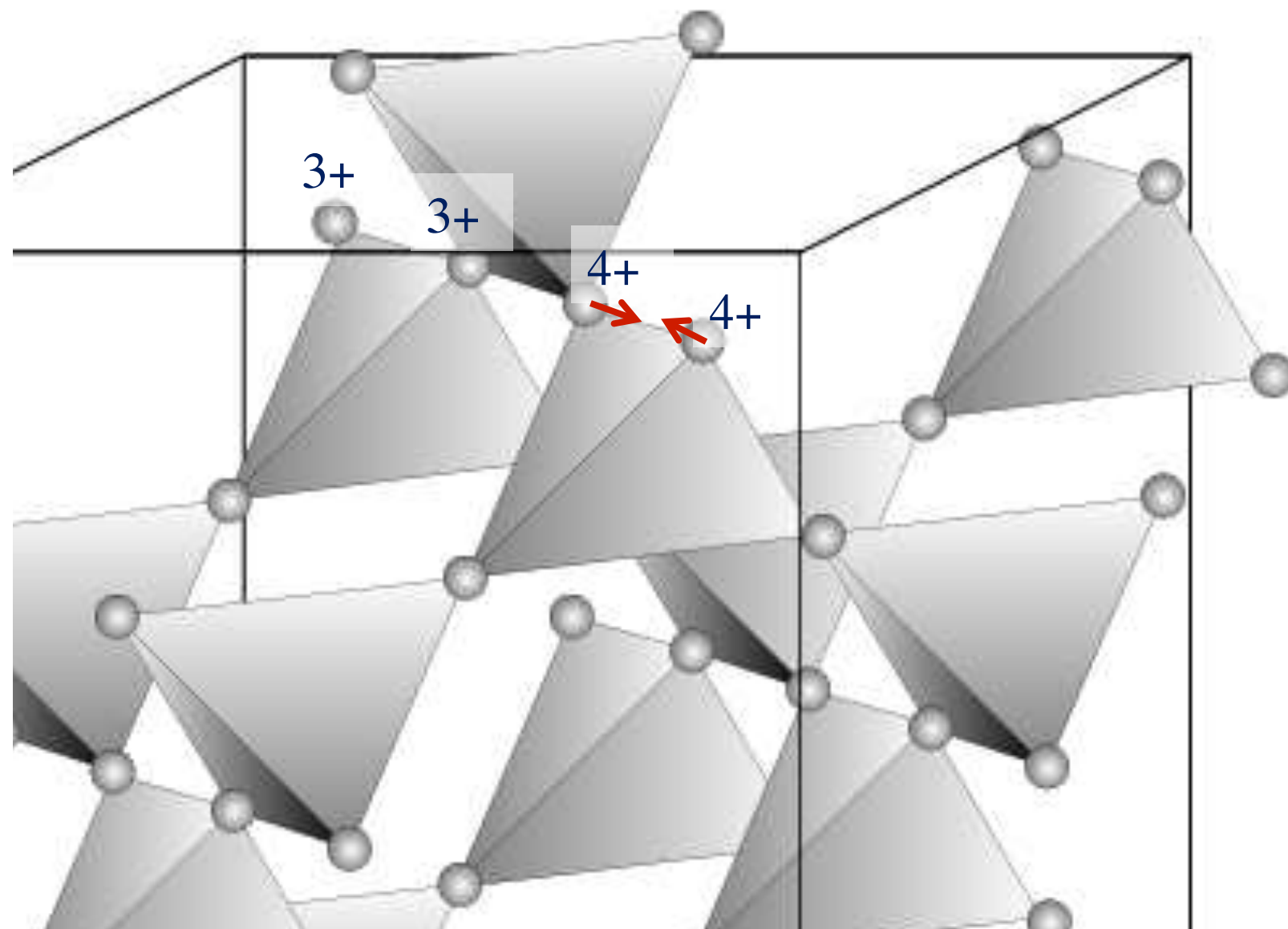
But the B atoms form a pyrochlore lattice of corner shared tetrahedra => low dimensional physics!





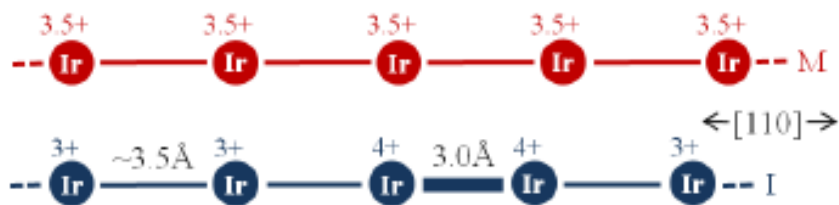
05).

[apam.columbia.edu/bgsite](http://apam.columbia.edu/bgsite)

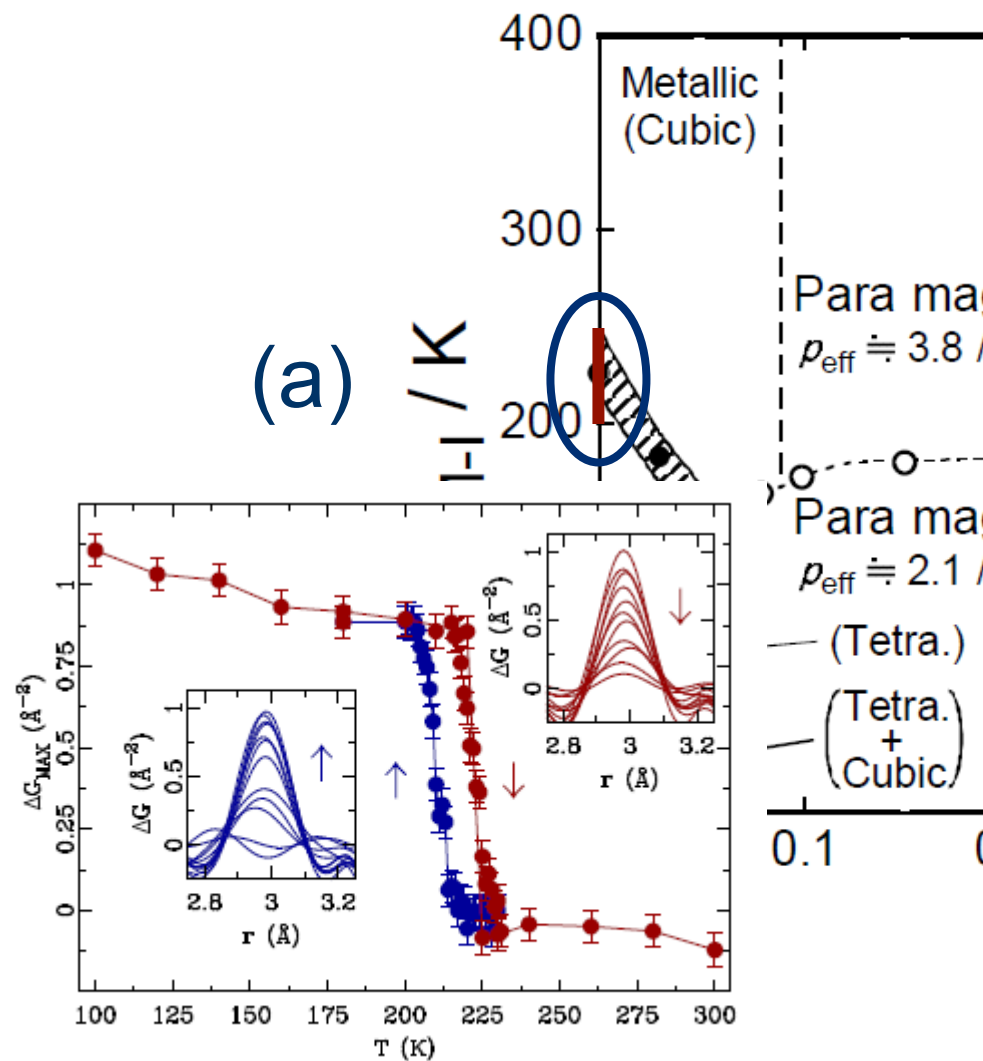
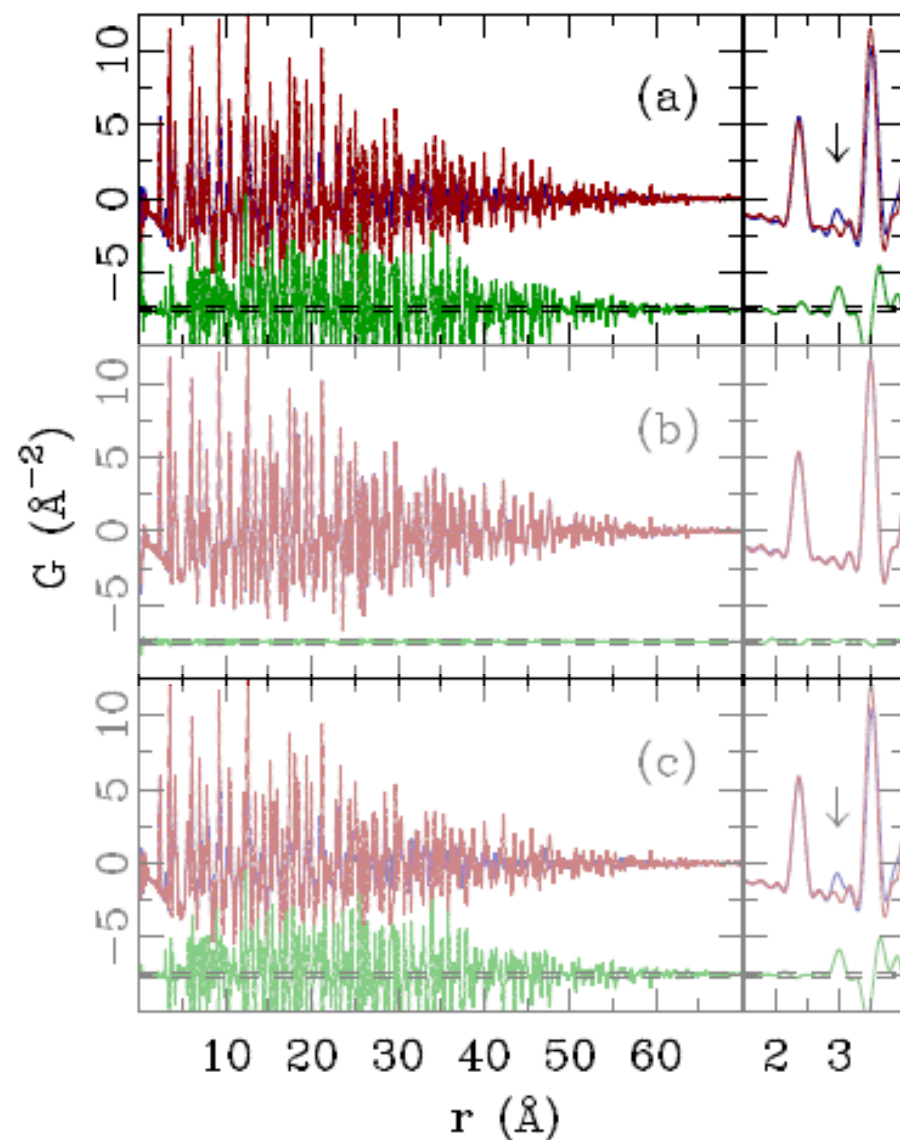


05).

[apam.columbia.edu/bgsite](http://apam.columbia.edu/bgsite)



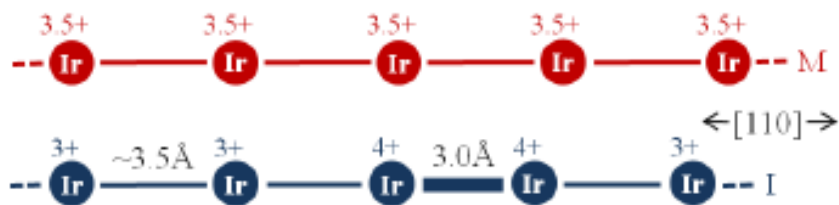
C



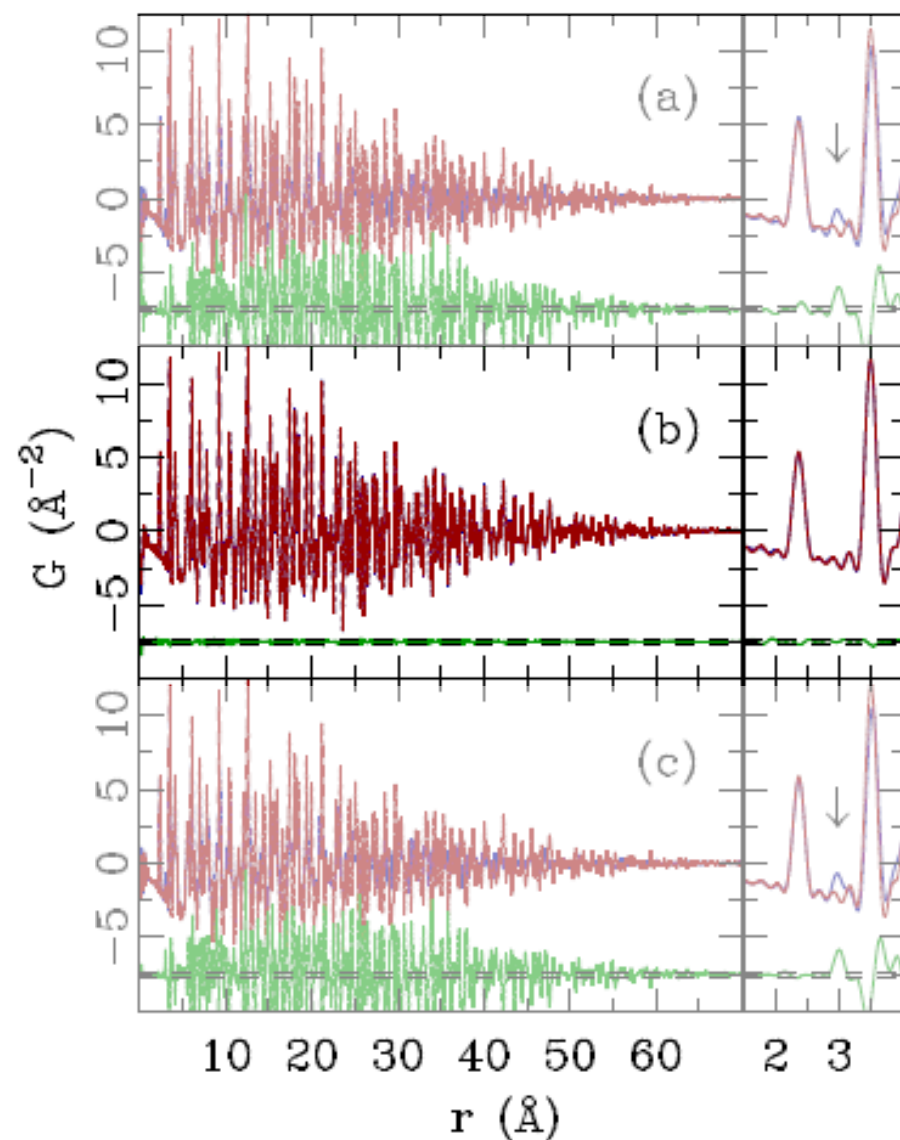
GENERATORY

[HTTP://bgsite.apam.columbia.edu/bgsite](http://bgsite.apam.columbia.edu/bgsite)

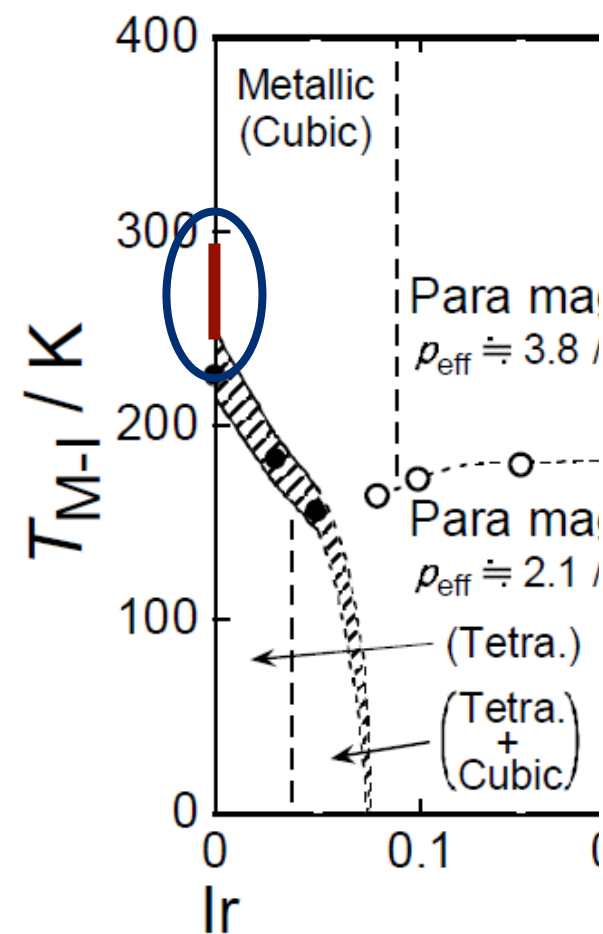




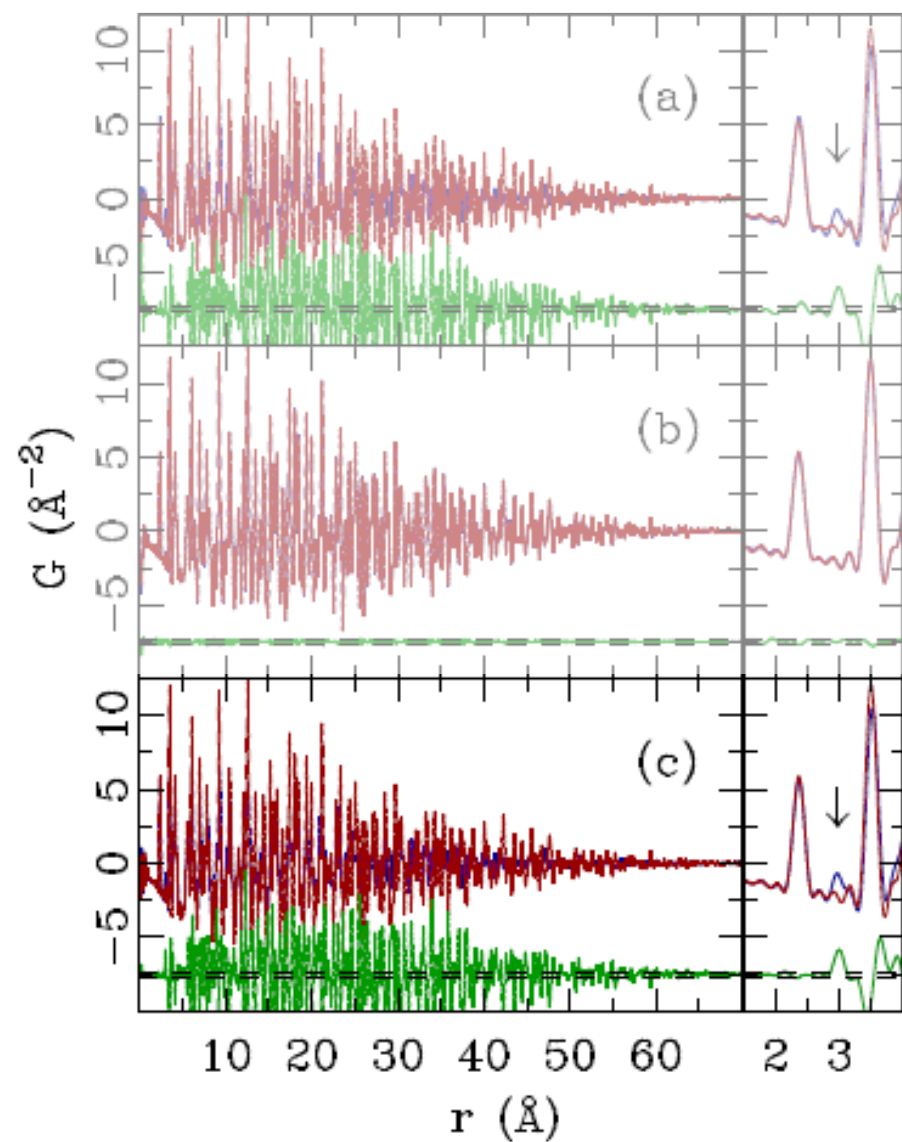
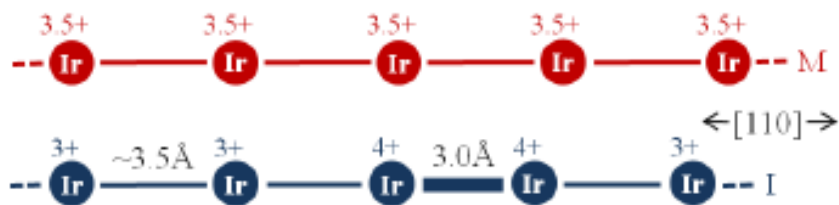
C



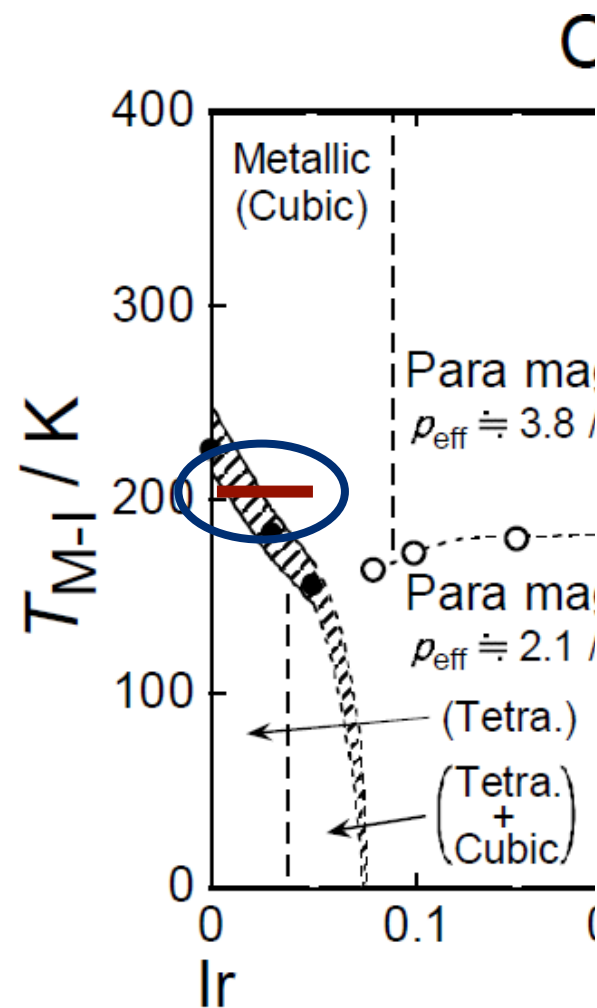
(b)





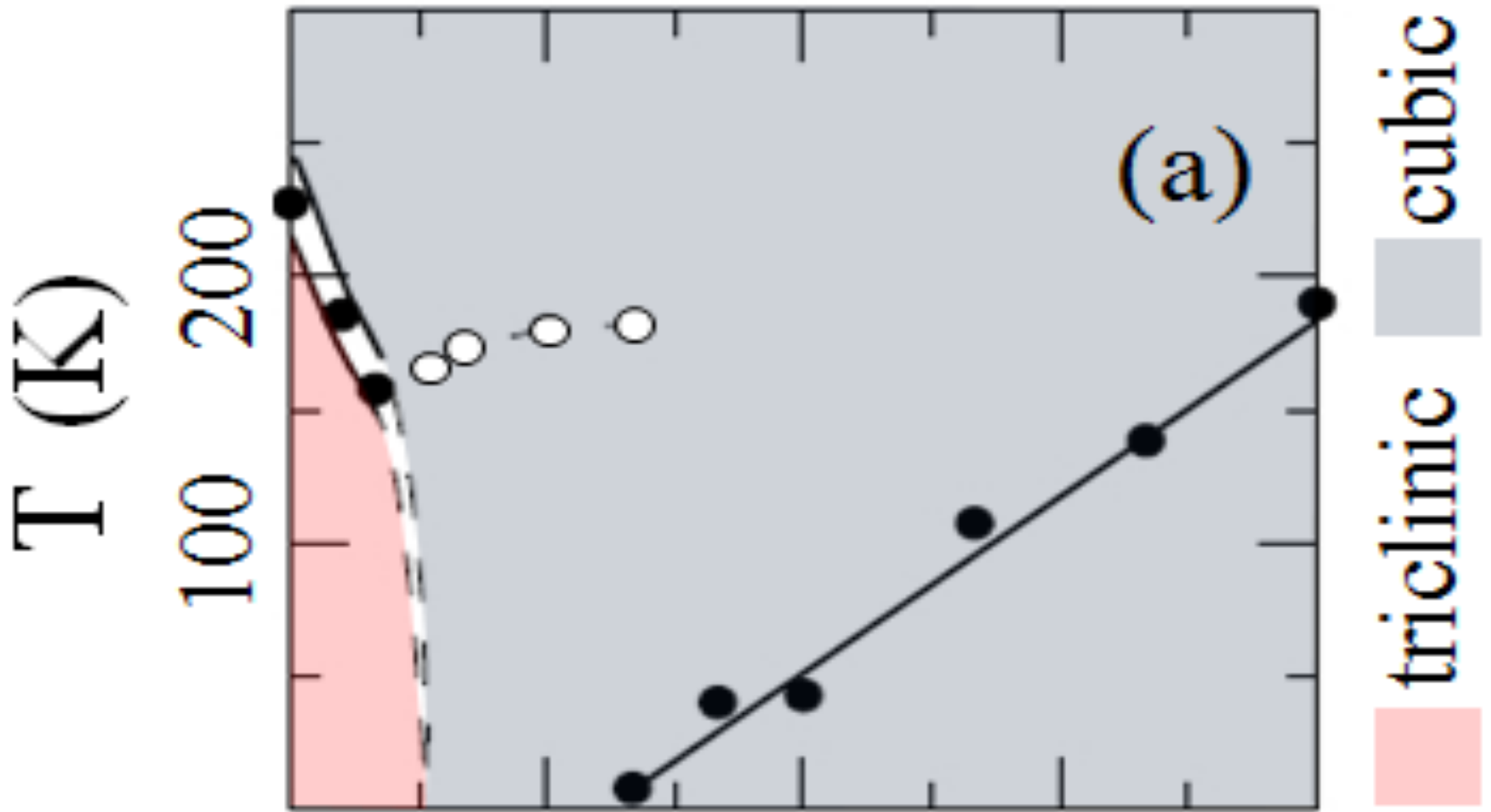


(c)

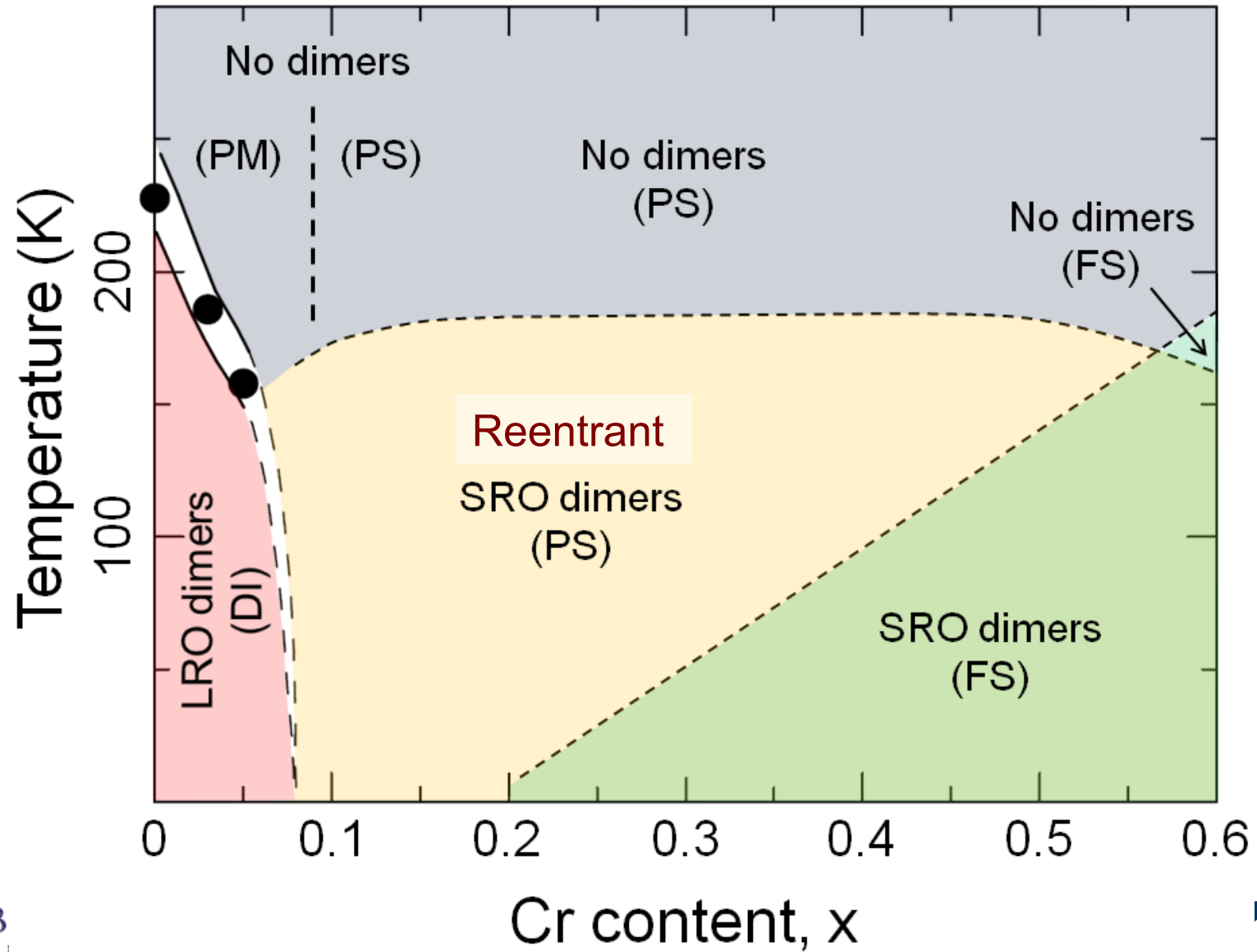




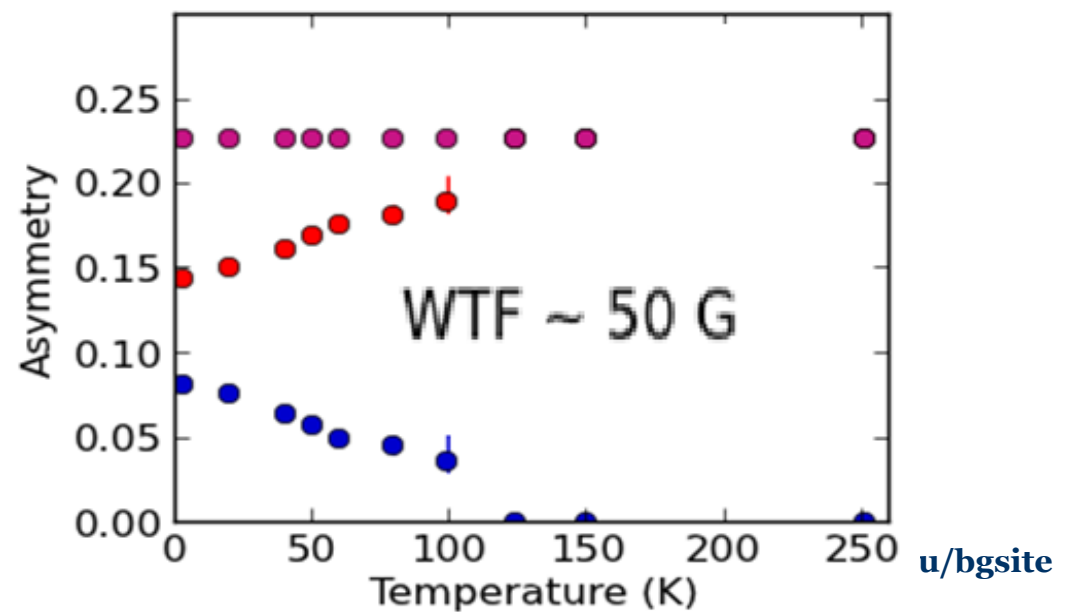
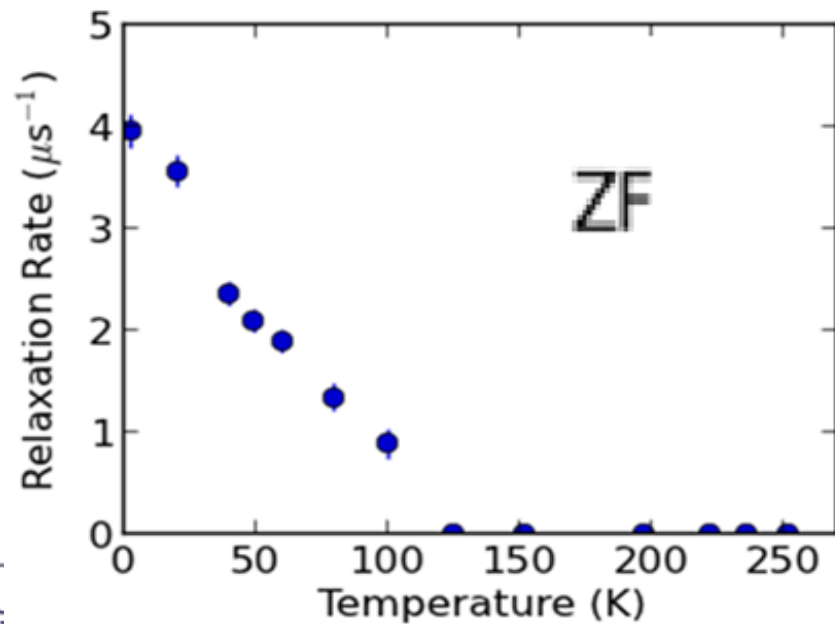
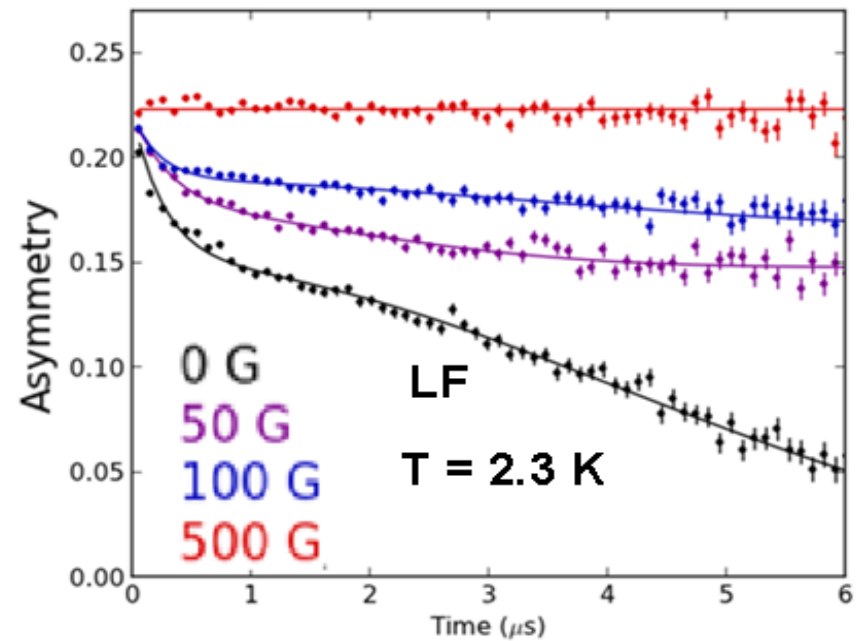
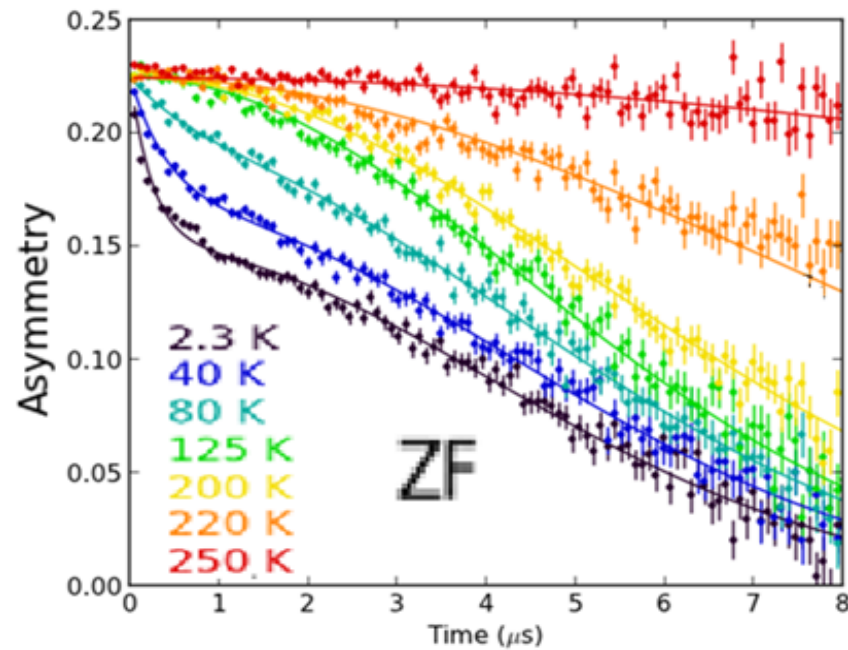
## Average phase diagram



## Local phase diagram



# Results: $x = 0$



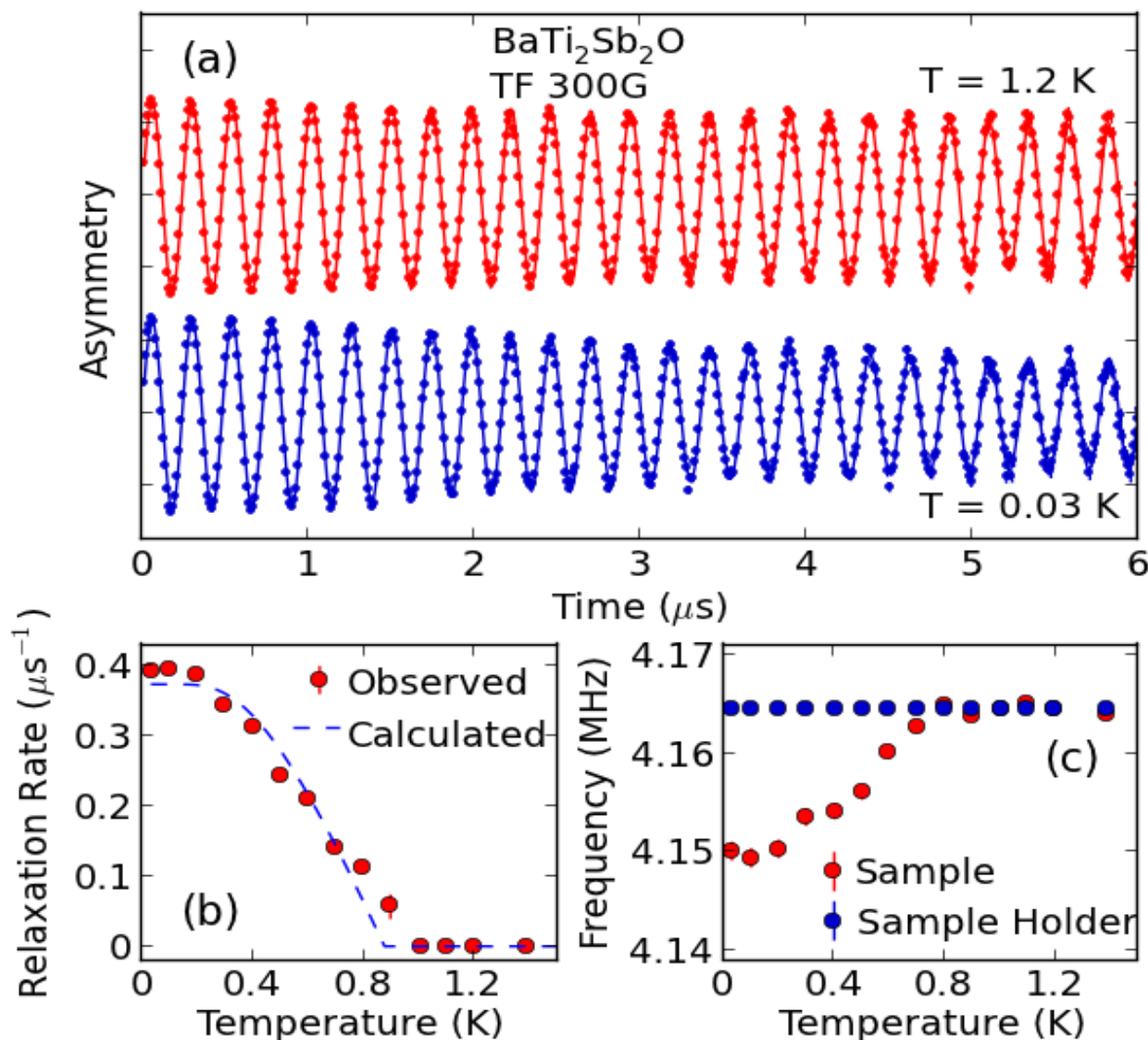


# Acknowledgements



- A special thank you to all my current and former students and post-docs
- Emil Bozin, Milinda Abeykoon, Pavol Juhas
- Also my many wonderful collaborators, in particular from the work described here: John Mitchell, John Tranquada, John Hill, Weiguo Yin, Genda Gu
- Facilities:
  - NSLS (and people therein)
  - Lujan Center LANL (and people therein)
  - Triumph, Canada
- Funding: DOE-BES

# Superconductivity probed by $\mu$ SR



Conclusion: fully gapped  
s-wave superconductivity

$T_c \approx 1\text{ K}$

$\lambda_L (T = 0) \approx 430\text{ nm}$

Agreement with other  $\mu$ SR  
and heat capacity studies  
(von Rohr et al, PRB 88:  
140501 (2013); Gooch et  
al, PRB 88: 064510 (2013))

Y. Nozaki et al. PRB 88 214506 (2013)

[HTTP://bgsite.apam.columbia.edu/bgsite](http://bgsite.apam.columbia.edu/bgsite)